Adv. Oper. Theory 2 (2017), no. 3, 318-333
http://doi.org/10.22034/aot.1703-1133
ISSN: 2538-225X (electronic)
http://aot-math.org

# $k$ TH-ORDER SLANT TOEPLITZ OPERATORS ON THE FOCK SPACE 

SHIVAM KUMAR SINGH ${ }^{1 *}$ and ANURADHA GUPTA ${ }^{2}$<br>Communicated by E. Ko


#### Abstract

The notion of slant Toeplitz operators $B_{\phi}$ and $k$ th-order slant Toeplitz operators $B_{\phi}^{k}$ on the Fock space is introduced and some of its properties are investigated. The Berezin transform of slant Toeplitz operator $B_{\phi}$ is also obtained. In addition, the commutativity of $k$ th-order slant Toeplitz operators with co-analytic and harmonic symbols is discussed.


## 1. Introduction

O. Toeplitz [16] in 1911 introduced the notion of Toeplitz operator $T_{\phi}$ for bounded measurable function $\phi$ with applications in prediction theory, wavelet analysis and differential equations. Later on Toeplitz operators on Hardy spaces have been studied extensively. In particular authors like Brown, Halmos [4] and Douglas [5] have done the remarkable study of this operators. Then, in the year 1995, Ho [8, 9, 10, 11], introduced slant Toeplitz operator having the property that its matrix with respect to standard orthonormal basis could be obtained by eliminating every alternate row of the matrix of the corresponding Toeplitz operator. After the introduction of class of slant Toeplitz operators, the study has gained voluminous importance due to its multidirectional applications as these class of operators have played major roles in wavelet analysis, dynamical system and in curve and surface modelling ( $[6,7,13,14,17]$ ). Many mathematicians (for e.g. $[1,2,3])$ generalized the notion of slant Toeplitz operators to different spaces such as Hardy spaces, Bergman space and studied its properties. Motivated by

[^0]the work of these researchers, here we introduce slant Toeplitz operators on the Fock space and study some of its algebraic properties.

Let the space $L^{2}(\mathbb{C}, d \mu)$ denote the Hilbert space of all Lebesgue measurable square integrable functions $f$ on $\mathbb{C}$ with the norm

$$
\|f\|=\left(\int_{\mathbb{C}}|f(z)|^{2} d \mu(z)\right)^{1 / 2}
$$

where the measure $d \mu(z)=e^{-|z|^{2}} d A(z)$ and $d A$ denotes the Lebesgue area measure on complex plane. We denote the Fock space by $\mathbb{F}^{2}$ which consists of all entire functions in $L^{2}(\mathbb{C}, d \mu)$ and is a closed subspace of $L^{2}(\mathbb{C}, d \mu)$. The space $\mathbb{F}^{2}$ is a Hilbert space with the inner product inherited from $L^{2}(\mathbb{C}, d \mu)$ as

$$
\begin{equation*}
\langle f, g\rangle=\int_{\mathbb{C}} f(z) \overline{g(z)} d \mu(z) \quad \text { where } f \text { and } g \text { are in } \mathbb{F}^{2} . \tag{1.1}
\end{equation*}
$$

The set of all polynomials in the complex variable $z$ is denoted by $\mathbb{P}[z]$ which is contained in the Fock space and moreover it is dense. For $n \geq 0$, let $e_{n}(z)=$ $\frac{z^{n}}{\sqrt{\pi n!}}$, then the set $\left\{e_{n}\right\}_{n \geq 0}$ forms the orthonormal basis for $\mathbb{F}^{2}$ (see [18]). Let $P: L^{2}(\mathbb{C}, d \mu) \longrightarrow \mathbb{F}^{2}$ be the orthogonal projection. Then for $f \in L^{2}(\mathbb{C}, d \mu)$, we have

$$
\begin{aligned}
P(f(z)) & =\left\langle P f(w), K_{z}(w)\right\rangle=\left\langle f(w), P\left(K_{z}(w)\right)\right\rangle \\
& =\left\langle f(w), K_{z}(w)\right\rangle=\frac{1}{\pi} \int_{\mathbb{C}} f(w) e^{z \bar{w}} d \mu(w)
\end{aligned}
$$

where $K_{z}(w)=\overline{K(z, w)}=\frac{1}{\pi} e^{w \bar{z}}$ is the Fock kernel and the normalized reproducing kernel is given by

$$
k_{z}(w)=\frac{K_{z}(w)}{\left\|K_{z}\right\|}=\frac{K(w, z)}{\sqrt{K(z, z)}}=\frac{1}{\sqrt{\pi}} e^{w \bar{z}-\frac{|z|^{2}}{2}} .
$$

Let $L^{\infty}(\mathbb{C})$ denote the set of all essentially bounded measurable functions in the entire complex plane, then for $\phi \in L^{\infty}(\mathbb{C})$, the multiplication operator on $L^{2}(\mathbb{C}, d \mu)$ is defined by $M_{\phi}(f)=\phi \cdot f$ for all $f \in L^{2}(\mathbb{C}, d \mu)$ and the Toeplitz operator $T_{\phi}$ on $\mathbb{F}^{2}$ is defined by $T_{\phi}(f)=P(\phi \cdot f)$ for all $f \in \mathbb{F}^{2}$.

In this paper, we have introduced the notion of slant Toeplitz operators $B_{\phi}$ and the $k$ th-order slant Toeplitz operators $B_{\phi}^{k}$ on the Fock space and have studied its properties. In particular, we have given the explicit expression for Berezin transform of slant Toeplitz operator $B_{\phi}$ and also obtained the conditions for boundedness, compactness of $B_{\phi}$. In addition, we have shown that the necessary and sufficient conditions for the commutativity of $k$ th-order slant Toeplitz operators on the Fock space are that their symbols functions must be linearly dependent.

## 2. $k$ th-order slant Toeplitz operators on the Fock space

Lemma 2.1. For non-negative integers $s$ and $t$, we have

$$
\left\langle z^{s}, z^{t}\right\rangle= \begin{cases}\pi s! & \text { if } s=t \\ 0 & \text { otherwise }\end{cases}
$$

Proof. Indeed for $s=t$, using the equation (1.1) and the measure $d \mu(z)$, we get

$$
\begin{aligned}
\left\langle z^{s}, z^{t}\right\rangle & =\int_{\mathbb{C}} z^{s} z^{t} e^{-|z|^{2}} d A(z)=\int_{r=0}^{\infty} \int_{\theta=0}^{2 \pi} r^{s+t+1} e^{i(s-t) \theta} e^{-r^{2}} d \theta d r \\
& =2 \pi \int_{r=0}^{\infty} r^{2 s+1} e^{-r^{2}} d r=\pi \int_{l=0}^{\infty} l^{s} e^{-l} d l=\pi \Gamma(s+1)=\pi s!
\end{aligned}
$$

Then $\left\langle z^{s}, z^{t}\right\rangle=\left\{\begin{array}{ll}\pi s! & \text { if } s=t \\ 0 & \text { otherwise }\end{array}\right.$.
Now in the next proposition we give the $(m, n)^{\text {th }}$ entry of the matrix of $T_{\phi}$ with respect to orthonormal basis $\left\{e_{n}\right\}_{n \geq 0}$ on the Fock space.

Proposition 2.2. For the harmonic symbol $\phi(z)=\sum_{i=0}^{\infty} a_{i} z^{i}+\sum_{j=1}^{\infty} b_{j} \bar{z}^{j}$, the $(m, n)^{\text {th }}$ entry of matrix of $T_{\phi}$ with respect to orthonormal basis $\left\{e_{n}\right\}_{n \geq 0}$ of $\mathbb{F}^{2}$ is given by

$$
\left\langle T_{\phi} e_{n}, e_{m}\right\rangle= \begin{cases}\sqrt{\frac{m!}{n!}} a_{m-n} & \text { for } m \geq n \\ \sqrt{\frac{n!}{m!}} b_{n-m} & \text { for } n>m\end{cases}
$$

where $m$ and $n$ are non-negative integers.
Proof. Here the $(m, n)^{t h}$ entry of matrix of $T_{\phi}$ with respect to orthonormal basis $\left\{e_{n}\right\}_{n \geq 0}$ of $\mathbb{F}^{2}$ is given by

$$
\begin{align*}
\left\langle T_{\phi} e_{n}, e_{m}\right\rangle & =\frac{1}{\sqrt{\pi n!}} \frac{1}{\sqrt{\pi m!}}\left\langle T_{\phi} z^{n}, z^{m}\right\rangle \\
& =\frac{1}{\pi \sqrt{n!}} \frac{1}{\sqrt{m!}}\left\langle P\left(\sum_{i=0}^{\infty} a_{i} z^{i+n}+\sum_{j=1}^{\infty} b_{j} \bar{z}^{j} z^{n}\right), z^{m}\right\rangle \\
& =\frac{1}{\pi \sqrt{n!}} \frac{1}{\sqrt{m!}}\left(\left\langle\sum_{i=0}^{\infty} a_{i} z^{i+n}, z^{m}\right\rangle+\left\langle\sum_{j=1}^{\infty} b_{j} \bar{z}^{j} z^{n}, z^{m}\right\rangle\right) \\
& =\frac{1}{\pi \sqrt{n!}} \frac{1}{\sqrt{m!}}\left(\sum_{i=0}^{\infty} a_{i}\left\langle z^{i+n}, z^{m}\right\rangle+\sum_{j=1}^{\infty} b_{j}\left\langle z^{n}, z^{j+m}\right\rangle\right) \tag{2.1}
\end{align*}
$$

For $m \geq n$, by Lemma 2.1 and by equation (2.1), it follows that

$$
\begin{aligned}
\left\langle T_{\phi} e_{n}, e_{m}\right\rangle & =\frac{1}{\sqrt{\pi n!}} \frac{1}{\sqrt{\pi m!}} \sum_{i=0}^{\infty} a_{i}\left\langle z^{i+n}, z^{m}\right\rangle \\
& =\frac{1}{\sqrt{\pi n!}} \frac{1}{\sqrt{\pi m!}} a_{m-n} \pi m!=\sqrt{\frac{m!}{n!}} a_{m-n}
\end{aligned}
$$

For $n>m$, by Lemma 2.1 and by equation (2.1), it follows that

$$
\begin{gathered}
\qquad \begin{array}{c}
\left\langle T_{\phi} e_{n}, e_{m}\right\rangle=\frac{1}{\sqrt{\pi n!}} \frac{1}{\sqrt{\pi m!}} \sum_{j=1}^{\infty} b_{j}\left\langle z^{n}, z^{j+m}\right\rangle \\
=\frac{1}{\sqrt{\pi n!}} \frac{1}{\sqrt{\pi m!}} b_{n-m} \pi(n)!=\sqrt{\frac{n!}{m!}} b_{n-m} \\
\text { Thus }\left\langle T_{\phi} e_{n}, e_{m}\right\rangle= \begin{cases}\sqrt{\frac{m!}{n!}} a_{m-n} & \text { for } m \geq n \\
\sqrt{\frac{n!}{m!}} b_{n-m} & \text { for } n>m\end{cases}
\end{array} .\left\{\begin{array}{l}
\end{array}\right.
\end{gathered}
$$

where $m$ and $n$ are non-negative integers.
Hence the matrix of $T_{\phi}$ explicitly is given by

$$
T_{\phi}=\left(\begin{array}{ccccc}
a_{0} & b_{1} & \sqrt{2} b_{2} & \sqrt{6} b_{3} & \ldots \\
a_{1} & a_{0} & \sqrt{2} b_{1} & \sqrt{6} b_{2} & \ldots \\
\sqrt{2} a_{2} & \sqrt{2} a_{1} & a_{0} & \sqrt{3} b_{1} & \ldots \\
\sqrt{6} a_{3} & \sqrt{6} a_{2} & \sqrt{3} a_{1} & a_{0} & \cdots \\
\vdots & \vdots & \vdots & \vdots &
\end{array}\right)
$$

and the adjoint of the matrix of $T_{\phi}$ is given by

$$
T_{\phi}^{\star}=\left(\begin{array}{ccccc}
\overline{a_{0}} & \overline{a_{1}} & \sqrt{2} \overline{a_{2}} & \sqrt{6} \overline{a_{3}} & \ldots \\
\overline{b_{1}} & \overline{a_{0}} & \sqrt{2} \overline{a_{1}} & \sqrt{6} \overline{a_{2}} & \ldots \\
\sqrt{2} \overline{b_{2}} & \sqrt{2} \overline{b_{1}} & \overline{a_{0}} & \sqrt{3} \overline{a_{1}} & \ldots \\
\sqrt{6} \overline{b_{3}} & \sqrt{6} \overline{b_{2}} & \sqrt{3} \overline{b_{1}} & \overline{a_{0}} & \ldots \\
\vdots & \vdots & \vdots & \vdots &
\end{array}\right)
$$

which is the matrix of $T_{\hat{\phi}}$ where $\hat{\phi}(z)=\sum_{i=0}^{\infty} \overline{a_{i}} \bar{z}^{i}+\sum_{j=1}^{\infty} \bar{b}_{j} z^{j}$ which is nothing but $\bar{\phi}(z)$. So we get that $T_{\phi}^{\star}=T_{\bar{\phi}}$.
Assume that $k$ is an integer and $k \geq 2$ and now consider an operator $W_{k}$ on $\mathbb{F}^{2}$ given by

$$
W_{k}\left(z^{n}\right)=\left\{\begin{array}{ll}
z^{\frac{n}{k}}, & \text { if } n \text { is divisible by } k \\
0 & \text { otherwise }
\end{array} .\right.
$$

Note here we denote the operator $W_{2}$ by $W$.
Proposition 2.3. The operator $W_{k}$ is a bounded linear operator on $\mathbb{F}^{2}$ with norm 1.

Proof. Let $f(z)=\sum_{n=0}^{\infty} a_{n} z^{n}$ be any arbitrary polynomial, then

$$
\begin{aligned}
\left\|W_{k} f\right\|_{2}^{2} & =\left\langle\sum_{m=0}^{\infty} a_{k m} z^{m}, \sum_{m=0}^{\infty} a_{k m} z^{m}\right\rangle=\sum_{m=0}^{\infty}\left|a_{k m}\right|^{2} \pi m! \\
& \leq \pi \sum_{m=0}^{\infty}\left|a_{m}\right|^{2} m!=\|f\|_{2}^{2}
\end{aligned}
$$

and since the set of all polynomials is dense in $\mathbb{F}^{2}$, so it follows that $\left\|W_{k}\right\|_{2} \leq 1$. Also $\left\|W_{k}\right\|_{2} \geq\left\|W_{k}\left(\frac{1}{\sqrt{\pi}}\right)\right\|_{2}=1$. Hence $\left\|W_{k}\right\|_{2}=1$.
Proposition 2.4. For the operator $W_{k}$, its adjoint $W_{k}^{\star}\left(z^{n}\right)=\frac{n!}{(k n)!} z^{k n}$ for $n=0,1,2,3 \cdots$ and also $\left\|W_{k}^{\star} f\right\|_{2} \leq\|f\|_{2}$, for all $f \in \mathbb{F}^{2}$.
Proof. Since the set of polynomials is dense in $\mathbb{F}^{2}$, therefore for any polynomial $f(z)=\sum_{m=0}^{\infty} a_{m} z^{m} \in \mathbb{F}^{2}$, we have

$$
\begin{aligned}
\left\langle W_{k}^{\star} z^{n}, f(z)\right\rangle & =\left\langle z^{n}, W_{k}\left(\sum_{m=0}^{\infty} a_{m} z^{m}\right)\right\rangle=\left\langle z^{n}, \sum_{m=0}^{\infty} a_{k m} z^{m}\right\rangle \\
& =\left\langle z^{n}, a_{k n} z^{n}\right\rangle=\overline{a_{k n}} \pi n!=\overline{a_{k n}} \frac{\pi n!}{\pi(k n)!}\left\langle z^{k n}, z^{k n}\right\rangle \\
& =\left\langle\frac{n!}{(k n)!} z^{k n}, f(z)\right\rangle
\end{aligned}
$$

This implies that $W_{k}^{\star}\left(z^{n}\right)=\frac{n!}{(k n)!} z^{k n}$. Since

$$
\begin{aligned}
\left\|W_{k}^{\star} f\right\|_{2}^{2} & =\left\|W_{k}^{\star} \sum_{m=0}^{\infty} a_{m} z^{m}\right\|_{2}^{2}=\left\|\sum_{m=0}^{\infty} a_{m} \frac{m!}{(k m)!} z^{k m}\right\|_{2}^{2} \\
& =\sum_{m=0}^{\infty}\left|a_{m}\right|^{2}\left(\frac{m!}{(k m)!}\right)^{2}\left\langle z^{k m}, z^{k m}\right\rangle=\sum_{m=0}^{\infty}\left|a_{m}\right|^{2}\left(\frac{m!}{(k m)!}\right)^{2} \pi(k m)! \\
& =\pi \sum_{m=0}^{\infty}\left|a_{m}\right|^{2} \frac{m!^{2}}{(k m)!} \leq \pi \sum_{m=0}^{\infty}\left|a_{m}\right|^{2} m!=\|f\|_{2}^{2}
\end{aligned}
$$

We have $\left\|W_{k}^{\star} f\right\|_{2} \leq\|f\|_{2}$ for all $f \in \mathbb{F}^{2}$.
Now we define slant Toeplitz operator and $k^{t h}$-order slant Toeplitz operator on the Fock space.
Definition 2.5. For $\phi \in L^{\infty}$, we define the slant Toeplitz operator on the space $\mathbb{F}^{2}$ as an operator

$$
B_{\phi}: \mathbb{F}^{2} \longrightarrow \mathbb{F}^{2} \quad \text { given by } \quad B_{\phi}(f)=W T_{\phi}(f) \quad \forall f \in \mathbb{F}^{2}
$$

Definition 2.6. A $k^{\text {th }}$-order slant Toeplitz operator $B_{\phi}^{k}$ for $k \geq 2$ induced by a function $\phi \in L^{\infty}$ on the space $\mathbb{F}^{2}$ is defined as

$$
B_{\phi}^{k}=W_{k} T_{\phi}
$$

It is clear that $B_{\phi}^{k}$ is bounded linear operator on $\mathbb{F}^{2}$ for essentially bounded measurable function $\phi$ over $\mathbb{C}$ and for $k=2$, the $k^{t h}$-order slant Toeplitz operator is simply the slant Toeplitz operator $B_{\phi}$.
Following proposition follows easily from the definition of $k^{t h}$-order slant Toeplitz operator $B_{\phi}^{k}$.
Proposition 2.7. Let $\phi_{1}, \phi_{2} \in L^{\infty}(\mathbb{C})$ and $\lambda_{1}, \lambda_{2}$ be complex numbers. Then
(1) $B_{\lambda_{1} \phi_{1}+\lambda_{2} \phi_{2}}^{k}=\lambda_{1} B_{\phi_{1}}^{k}+\lambda_{2} B_{\phi_{2}}^{k}$.

$$
\text { (2) }\left\|B_{\phi_{1}}^{k}\right\| \leq\left\|\phi_{1}\right\|_{\infty} \text {. }
$$

Proposition 2.8. For $\phi(z)=\sum_{i=0}^{\infty} a_{i} z^{i}+\sum_{j=1}^{\infty} b_{j} \bar{z}^{j}$, the ( $\left.m, n\right)^{\text {th }}$ entry of the matrix of $B_{\phi}^{k}$ with respect to orthonormal basis $\left\{e_{n}\right\}_{n \geq 0}$ of $\mathbb{F}^{2}$ is given by

$$
\left\langle B_{\phi}^{k} e_{n}, e_{m}\right\rangle= \begin{cases}\sqrt{\frac{m!}{n!}} a_{k m-n} & \text { for } k m \geq n \\ \frac{\sqrt{(m)!(n)!}}{(k m)!} b_{n-k m} & \text { for } n>k m\end{cases}
$$

where $m$ and $n$ are non-negative integers.
Proof. Here the $(m, n)^{t h}$ entry of $B_{\phi}^{k}$ with respect to orthonormal basis $\left\{e_{n}\right\}_{n \geq 0}$ of $\mathbb{F}^{2}$ is

$$
\begin{align*}
\left\langle B_{\phi}^{k} e_{n}, e_{m}\right\rangle & =\left\langle W_{k} T_{\phi} \frac{z^{n}}{\sqrt{\pi n!}}, \frac{z^{m}}{\sqrt{\pi m!}}\right\rangle \\
& =\frac{1}{\pi \sqrt{n!m!}}\left\langle P\left(\phi \cdot z^{n}\right), W_{k}^{\star} z^{m}\right\rangle \\
& =\frac{1}{\pi \sqrt{n!m!}}\left\langle\sum_{i=0}^{\infty} a_{i} z^{i+n}+\sum_{j=1}^{\infty} b_{j} \bar{z}^{j} z^{n}, \frac{m!}{(k m)!} z^{k m}\right\rangle \\
& =\frac{1}{\pi \sqrt{n!m!}}\left(\sum_{i=0}^{\infty} a_{i}\left\langle z^{i+n}, \frac{m!}{(k m)!} z^{k m}\right\rangle+\sum_{j=1}^{\infty} b_{j}\left\langle\bar{z}^{j} z^{n}, \frac{m!}{(k m)!} z^{k m}\right\rangle\right) \\
& =\frac{1}{\pi(k m)!} \sqrt{\frac{m!}{n!}}\left(\sum_{i=0}^{\infty} a_{i}\left\langle z^{i+n}, z^{k m}\right\rangle+\sum_{j=1}^{\infty} b_{j}\left\langle z^{n}, z^{k m+j}\right\rangle\right) . \tag{2.2}
\end{align*}
$$

For $k m \geq n$, by Lemma 2.1 and by equation (2.2), it follows that

$$
\begin{aligned}
\left\langle B_{\phi}^{k} e_{n}, e_{m}\right\rangle & =\frac{1}{\pi(k m)!} \sqrt{\frac{m!}{n!}} \sum_{i=0}^{\infty} a_{i}\left\langle z^{i+n}, z^{k m}\right\rangle \\
& =\frac{1}{\pi} \frac{1}{(k m)!} \sqrt{\frac{m!}{n!}} a_{k m-n} \pi(k m)!=\sqrt{\frac{m!}{n!}} a_{k m-n} .
\end{aligned}
$$

For $n>k m$, by Lemma 2.1 and by equation (2.2), it follows that

$$
\begin{aligned}
\left\langle B_{\phi}^{k} e_{n}, e_{m}\right\rangle & =\frac{1}{\pi(k m)!} \sqrt{\frac{m!}{n!}} \sum_{j=1}^{\infty} b_{j}\left\langle z^{n}, z^{k m+j}\right\rangle \\
& =\frac{1}{\pi(k m)!} \sqrt{\frac{m!}{n!}} b_{n-k m} \pi(n)!=\frac{1}{(k m)!} \sqrt{m!n!} b_{n-k m}
\end{aligned}
$$

Thus,

$$
\left\langle B_{\phi}^{k} e_{n}, e_{m}\right\rangle= \begin{cases}\sqrt{\frac{m!}{n!}} a_{k m-n} & \text { for } k m \geq n \\ \frac{\sqrt{(m)!(n)!}}{(k m)!} b_{n-k m} & \text { for } n>k m\end{cases}
$$

where $m$ and $n$ are non-negative integers.

Hence explicit form of the matrix of $B_{\phi}^{k}$ is given by

$$
\left(\begin{array}{cccc}
a_{0} & b_{1} & \sqrt{2} b_{2} & \ldots \\
a_{k} & a_{k-1} & \frac{1}{\sqrt{2}} a_{k-2} & \ldots \\
\sqrt{2} a_{2 k} & \sqrt{2} a_{2 k-1} & a_{2 k-2} & \ldots \\
\sqrt{6} a_{3 k} & \sqrt{6} a_{3 k-1} & \sqrt{3} a_{3 k-2} & \ldots \\
\vdots & \vdots & \vdots & \ldots
\end{array}\right)
$$

For $k=2$ we get the matrix of slant Toeplitz operator $B_{\phi}$, which is given by

$$
\left(\begin{array}{cccc}
a_{0} & b_{1} & \sqrt{2} b_{2} & \ldots \\
a_{2} & a_{1} & \frac{1}{\sqrt{2}} a_{0} & \ldots \\
\sqrt{2} a_{4} & \sqrt{2} a_{3} & a_{2} & \ldots \\
\sqrt{6} a_{6} & \sqrt{6} a_{5} & \sqrt{3} a_{4} & \ldots \\
\vdots & \vdots & \vdots & \ldots
\end{array}\right)
$$

Following proposition follows from the matrix representation of $B_{\phi}^{k}$.
Proposition 2.9. Let $\phi \in L^{\infty}$ be analytic function then $B_{\phi}^{k}=0$ if and only if $\phi=0$ almost everywhere.

Remark 2.10. For an analytic function $\phi \in L^{\infty}$ the correspondence $\phi \longrightarrow B_{\phi}^{k}$ is one-one.

It is evident that if $T_{\phi}$ is bounded on $\mathbb{F}^{2}$ then $B_{\phi}$ is always bounded. However reverse statement need not be true which is shown in the following example.

Example 2.11. $B_{z}$ and $B_{\bar{z}}$ are bounded linear operators on $\mathbb{F}^{2}$ while $T_{z}$ and $T_{\bar{z}}$ are not bounded on $\mathbb{F}^{2}$.

For any non-negtive integer $p$, using the lemma 2.1 we have

$$
\begin{aligned}
& \left\|B_{z} \cdot z^{p}\right\|^{2}=\left\|W T_{z} \cdot z^{p}\right\|^{2}=\left\|W z^{p+1}\right\|^{2} \\
& = \begin{cases}\left\|z^{\frac{p+1}{2}}\right\|^{2}, & \text { if } p \text { is odd } \\
0, & \text { otherwise }\end{cases} \\
& = \begin{cases}\frac{\left(\frac{p+1}{2}\right)!}{p!}\left\|z^{p}\right\|^{2}, & \text { if } p \text { is odd } \\
0, & \text { otherwise }\end{cases}
\end{aligned}
$$

which shows $B_{z}$ is a bounded linear operator.

Now for $p \in \mathbb{N}$, consider

$$
\begin{aligned}
& \left\|B_{\bar{z}} \cdot z^{p}\right\|^{2}=\left\|W T_{\bar{z}} \cdot z^{p}\right\|^{2}=\| W P\left(\bar{z} \cdot z^{p} \|^{2}\right. \\
& =\left\|W \frac{p!}{(p-1)!} z^{p-1}\right\|^{2}= \begin{cases}p^{2}\left\|z^{\frac{p-1}{2}}\right\|^{2}, & \text { if } p \text { is odd } \\
0, & \text { otherwise }\end{cases} \\
& = \begin{cases}p \frac{\left(\frac{p-1}{2}\right)!}{(p-1)!}\left\|z^{p}\right\|^{2}, & \text { if } p \text { is odd } \\
0, & \text { otherwise }\end{cases}
\end{aligned}
$$

which implies that $B_{\bar{z}}$ is bounded. Now consider

$$
\left\|T_{z} z^{p}\right\|^{2}=\left\langle z^{p+1}, z^{p+1}\right\rangle=\pi(p+1)!=(p+1)\left\langle z^{p}, z^{p}\right\rangle=(p+1)\left\|z^{p}\right\|^{2} .
$$

This means $\left\|T_{z} z^{p}\right\|=\sqrt{p+1}\left\|z^{p}\right\|$, thus $T_{z}$ is not bounded. Similarly we have $\left\|T_{\bar{z}} z^{p}\right\|=\sqrt{p}\left\|z^{p}\right\|$ and hence $T_{\bar{z}}$ is also not bounded.

Now we conclude this section by giving a result about the point spectra of $k$ th-order slant Toeplitz operators.

Theorem 2.12. Let $\phi$ be an invertible, analytic function in $L^{\infty}$, then $\sigma_{p}\left(B_{\phi}^{k}\right)=$ $\sigma_{p}\left(B_{\phi\left(z^{k}\right)}^{k}\right)$, where $\sigma_{p}\left(B_{\phi}^{k}\right)$ denotes the point spectrum of $B_{\phi}^{k}$.
Proof. Suppose that $\lambda \in \sigma_{p}\left(B_{\phi}^{k}\right)$. Then there exists a non-zero function $f$ in $\mathbb{F}^{2}$ such that $B_{\phi}^{k} f=\lambda f$. Let $M=\phi f$, then by using the definition of $W_{k}$ it follows that

$$
\begin{aligned}
B_{\phi}^{k}\left(z^{k}\right) M & =W_{k} T_{\phi\left(z^{k}\right)} M \\
& =W_{k} P\left[\phi\left(z^{k}\right)(\phi f)\right]=W_{k}\left[\phi\left(z^{k}\right)(\phi f)\right] \\
& =\phi(z) W_{k}(\phi f)=\phi(z) \cdot B_{\phi}^{k}(f) \\
& =\phi(z) \cdot \lambda f=\lambda \phi f \\
& =\lambda M .
\end{aligned}
$$

Since $\phi$ is invertible and $f$ is non-zero, so $M \neq 0$ and therefore $\lambda \in \sigma_{p}\left(B_{\phi}^{k}\right)$. Conversely, let $\mu \in \sigma_{p}\left(B_{\phi\left(z^{k}\right)}^{k}\right)$. Then there exist a non-zero function $g$ in $\mathbb{F}^{2}$ satisfying $B_{\phi}^{k}\left(z^{k}\right) g=\mu g$. Let $N=\phi^{-1} g$. Then $N \in \mathbb{F}^{2}$ satisfies

$$
\begin{aligned}
B_{\phi}^{k} N & =W_{k} T_{\phi}\left(\phi^{-1} g\right) \\
& =W_{k} P\left(\phi \cdot\left(\phi^{-1} g\right)\right)=W_{k} g \\
& =\phi^{-1} \phi W_{k} g \\
& =\phi^{-1} W_{k}\left(\phi\left(z^{k}\right) g\right) \\
& =\phi^{-1} B_{\phi\left(z^{k}\right)}^{k} g \\
& =\phi^{-1} \mu g=\mu \phi^{-1} g \\
& =\mu N
\end{aligned}
$$

Since $\phi$ is invertible and $g$ is non-zero, so $N \neq 0$ and therefore, $\mu \in \sigma_{p}\left(B_{\phi}^{k}\right)$. Thus $\sigma_{p}\left(B_{\phi}^{k}\right)=\sigma_{p}\left(B_{\phi\left(z^{k}\right)}^{k}\right)$.

Corollary 2.13. For an invertible, analytic function $\phi$ in $L^{\infty}, \sigma_{p}\left(B_{\phi}\right)=\sigma_{p}\left(B_{\phi\left(z^{2}\right)}\right)$, where $\sigma_{p}\left(B_{\phi}\right)$ denotes the point spectrum of $B_{\phi}$.

## 3. Berezin Transform of $B_{\phi}$

Let $\mathbb{H}$ be any reproducing kernel Hilbert space on an open subset $\Omega$ of $\mathbb{C}$. For a bounded operator $S$ on $\mathbb{H}$, the Berezin transform [15] denoted by $\tilde{S}$, is the complex valued function on $\Omega$

$$
\tilde{S}(z)=\left\langle S k_{z}, k_{z}\right\rangle \quad \text { for } \quad z \in \Omega
$$

For every bounded operator $S$ on $\mathbb{H}$, the Berezin transform $\tilde{S}$ is a bounded function on $\Omega$.
The normalized reproducing kernel in the Fock space is given by

$$
k_{z}(w)=\frac{1}{\sqrt{\pi}} e^{w \bar{z}-\frac{|z|^{2}}{2}}=\frac{1}{\sqrt{\pi}}\left(\sum_{n=0}^{\infty} \frac{(w \bar{z})^{n}}{n!}\right) e^{-\frac{|z|^{2}}{2}}
$$

The following proposition gives the Berezin transform of the operator $W_{k}$ :
Proposition 3.1. For the operator $W_{k}$, we have $W_{k}\left(k_{z}(w)\right)=\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} \sum_{m=0}^{\infty} \frac{(w \bar{z})^{m}}{(k m)!}$ and its Berezin transform is given by

$$
\tilde{W}_{k}(z)=\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} \sum_{m=0}^{\infty} \frac{|z|^{2 m}}{(k m)!}
$$

Proof. By the definitions of operator $W_{k}$ and normalized reproducing kernel $k_{z}$, it follows that

$$
\begin{aligned}
W_{k}\left(k_{z}(w)\right) & =W_{k}\left(\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}+w \bar{z}}\right)=\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} W_{k}\left(\sum_{n=0}^{\infty} \frac{(w \bar{z})^{n}}{n!}\right) \\
& =\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} \sum_{m=0}^{\infty} \frac{(w \bar{z})^{m}}{(k m)!}
\end{aligned}
$$

and

$$
\begin{aligned}
\tilde{W}_{k}(z) & =\left\langle W_{k} k_{z}, k_{z}\right\rangle=\left\langle\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} \sum_{m=0}^{\infty} \frac{(w \bar{z})^{m}}{(k m)!}, k_{z}(w)\right\rangle \\
& =\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} \sum_{m=0}^{\infty}\left\langle\frac{(w \bar{z})^{m}}{(k m)!}, k_{z}(w)\right\rangle \\
& =\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} \sum_{m=0}^{\infty} \frac{(z \bar{z})^{m}}{(k m)!}=\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} \sum_{m=0}^{\infty} \frac{|z|^{2 m}}{(k m)!} .
\end{aligned}
$$

The explicit expression for the Berezin transform of the slant Toeplitz operator is given in the next result:

Theorem 3.2. For an operator $W$ its adjoint $W^{\star}\left(k_{z}(w)\right)=\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} \cosh (w \bar{z})$ and the Berezin transform of slant Toeplitz operator $B_{\phi}$ is

$$
\tilde{B}_{\phi}(z)=\frac{1}{2 \pi} \int_{\mathbb{C}} \phi(w) e^{-|z-w|^{2}} d A(w)+\frac{1}{2 \pi} \int_{\mathbb{C}} \phi(w) e^{-|z-w|^{2}} e^{-2(z \bar{w})} d A(w) .
$$

Proof. By the proposition 2.4 and the definition of normalized reproducing kernel $k_{z}$, it follows that

$$
\begin{aligned}
W^{\star}\left(k_{z}(w)\right) & =\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} W^{\star}\left(\sum_{n=0}^{\infty} \frac{(w \bar{z})^{n}}{n!}\right) \\
& =\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} \sum_{n=0}^{\infty} \frac{(w \bar{z})^{2 n}}{(2 n!)} \\
& =\frac{1}{\sqrt{\pi}} e^{-\frac{|z|^{2}}{2}} \cosh (w \bar{z}) .
\end{aligned}
$$

Now the Berezin transform of $B_{\phi}$ is given by

$$
\begin{align*}
\tilde{B}_{\phi}(z) & =\left\langle B_{\phi} k_{z}, k_{z}\right\rangle \\
& =\left\langle W T_{\phi} k_{z}, k_{z}\right\rangle \\
& =\left\langle T_{\phi} k_{z}, W^{\star} k_{z}\right\rangle \\
& =\left\langle P\left(\phi k_{z}\right), W^{\star} k_{z}\right\rangle \\
& =\frac{1}{\sqrt{\pi}} \int_{\mathbb{C}} \phi(w) k_{z}(w) e^{-\frac{|z|^{2}}{2}} \overline{\cosh (w \bar{z})} e^{-|w|^{2}} d A(w) \\
& =\frac{1}{\pi} \int_{\mathbb{C}} \phi(w) e^{w \bar{z}-|z|^{2}-|w|^{2}} \cosh (z \bar{w}) d A(w) \\
& =\frac{1}{\pi} \int_{\mathbb{C}} \phi(w) e^{w \bar{z}-|z|^{2}-|w|^{2}}\left(\frac{e^{z \bar{w}}+e^{-(z \bar{w})}}{2}\right) d A(w) \\
& =\frac{1}{2 \pi} \int_{\mathbb{C}} \phi(w) e^{w \bar{z}+z \bar{w}-\left(|z|^{2}+|w|^{2}\right)} d A(w)+\frac{1}{2 \pi} \int_{\mathbb{C}} \phi(w) e^{w \bar{z}-z \bar{w}-\left(|z|^{2}+|w|^{2}\right)} d A(w) \\
& =\frac{1}{2 \pi} \int_{\mathbb{C}} \phi(w) e^{-|z-w|^{2}} d A(w)+\frac{1}{2 \pi} \int_{\mathbb{C}} \phi(w) e^{-|z-w|^{2}} e^{-2(z \bar{w})} d A(w) . \tag{3.1}
\end{align*}
$$

In the expression (3.1), the term $\frac{1}{\pi} \int_{\mathbb{C}} \phi(w) e^{-|z-w|^{2}} d A(w)$ is basically the Berezin transform of $T_{\phi}$ or simply the Berezin transform of $\phi$. Also from (3.1) it is clear that $\tilde{B}_{\phi}(z) \rightarrow 0$ as $|z| \rightarrow \infty$.
For $\phi \in L^{\infty}$, we observe that
(1) $\tilde{B}_{\phi}(z) \in L^{\infty}(\mathbb{C}, d \lambda)$.
(2) $\left\|\tilde{B}_{\phi}(z)\right\|_{\infty} \leq\|\phi\|_{\infty}$.
(3) $\tilde{B}_{\lambda_{1} \phi_{1}+\lambda_{2} \phi_{2}}=\lambda_{1} \tilde{B}_{\phi_{1}}+\lambda_{2} \tilde{B}_{\phi_{2}}$.

Corollary 3.3. Let $\phi$ be essentially bounded measurable function on $\mathbb{C}$. Then $B_{\phi}$ $=0$ if and only if $\tilde{B}_{\phi}=0$.

Proposition 3.4. If $\phi(z)$ is a non-negative bounded measurable function on $\mathbb{C}$ and if $|z| \rightarrow \infty$ then $B_{\phi}$ is compact.

Proof. If $|z| \rightarrow \infty$ then $\tilde{B}_{\phi} \rightarrow 0$ and so by the expression (3.1) it follows that $\int_{\mathbb{C}} \phi(w) e^{-|z-w|^{2}} d A(w) \rightarrow 0$ as $z \rightarrow \infty$, that is, $\tilde{T}_{\phi}(z) \rightarrow 0$ as $z \rightarrow \infty$. But we know that for a given $\phi, T_{\phi}$ is compact on $\mathbb{F}^{2}$ if and only if $\tilde{T}_{\phi} \rightarrow 0$ as $z \rightarrow \infty$, ([15], Proposition 5.3). So it follows that $T_{\phi}$ is compact. Hence $B_{\phi}$ is compact on $\mathbb{F}^{2}$ 。

## 4. Commutativity of $k$ th-order slant Toeplitz operators

In this section the commutativity of $k$ th-order slant Toeplitz operators with coanalytic symbols and harmonic symbols has been studied. The following lemma follows from [12].

Lemma 4.1. Let $f$ and $g$ be analytic functions in $L^{\infty}(\mathbb{C})$, both of which are not identically zero and let $q \geq 0$ be an integer. If $f W_{k}^{\star} g=g W_{k}^{\star} f$, then the following are equivalent;
(1) $f^{(i)}(0)=0$ for any integers $i$ with $0 \leq i \leq q$ and $f^{(q+1)}(0) \neq 0$
(2) $g^{(i)}(0)=0$ for any integers $i$ with $0 \leq i \leq q$ and $g^{(q+1)}(0) \neq 0$.

Now in the next theorem we obtain the necessary and sufficient condition for the commutativity of $k$ th-order slant Toeplitz operator with co-analytic symbols.
Theorem 4.2. Let $\phi, \psi \in L^{\infty}(\mathbb{C})$ be such that $\bar{\phi}, \bar{\psi}$ are analytic functions then the following statements are equivalent:
(1) $B_{\phi}^{k}$ and $B_{\psi}^{k}$ commute ;
(2) there exist scalers $\alpha$ and $\beta$, not both zero, such that $\alpha \phi+\beta \psi=0$.

Proof. Firstly suppose that (2) holds. Without loss of generality assume that $\alpha \neq 0$, then $\phi=\gamma \psi$ where $\gamma=-\beta / \alpha$. Then

$$
B_{\phi}^{k} B_{\psi}^{k}=W_{k} T_{\phi} W_{k} T_{\psi}=W_{k} T_{\gamma \psi} W_{k} T_{\psi}=\gamma W_{k} T_{\psi} W_{k} T_{\psi}=W_{k} T_{\psi} W T_{\gamma \psi}=B_{\psi}^{k} B_{\phi}^{k}
$$

Conversely suppose that $B_{\phi}^{k}$ and $B_{\psi}^{k}$ commutes. Therefore we get that $B_{\phi}^{k \star} B_{\psi}^{k \star}(1)=B_{\psi}^{k \star} B_{\phi}^{k \star}(1)$, or equivalently, $\bar{\phi} W_{k}^{\star} \bar{\psi}=\bar{\psi} W_{k}^{\star} \bar{\phi}$.
Now we have the following three cases.
Case I. If $\phi \equiv 0$ or $\psi \equiv 0$, then the result is obvious.
Case II. If $\bar{\phi}(0) \neq 0$ and $\bar{\psi}(0) \neq 0$. Since $\bar{\phi}, \bar{\psi}$ are analytic functions in $L^{\infty}(\mathbb{C})$, so let $\bar{\phi}(z)=\sum_{r!}^{\infty} a_{r=0} z^{r}$ and $\bar{\psi}(z)=\sum_{s=0}^{\infty} b_{s} z^{s}$, then $a_{0} \neq 0, b_{0} \neq 0$ and $W_{k}^{\star}(\bar{\phi}(z))=$ $\sum_{r=0}^{\infty} a_{r} \frac{r!}{(k r)!} z^{k r}$ and $W_{k}^{\star}(\bar{\psi}(z))=\sum_{s=0}^{\infty} b_{s} \frac{s!}{(k s)!} z^{k s}$. Also

$$
\bar{\psi} W_{k}^{\star}(\bar{\phi}(z))=\left(\sum_{s=0}^{\infty} b_{s} z^{s}\right)\left(\sum_{r=0}^{\infty} a_{r} \frac{r!}{(k r)!} z^{k r}\right)=\sum_{s=0}^{\infty} \sum_{r=0}^{\infty} \frac{r!}{(k r)!} b_{s} a_{r} z^{s+k r}
$$

and

$$
\bar{\phi} W_{k}^{\star}(\bar{\psi}(z))=\left(\sum_{r=0}^{\infty} a_{r} z^{r}\right)\left(\sum_{s=0}^{\infty} b_{s} \frac{s!}{(k s)!} z^{k s}\right)=\sum_{r=0}^{\infty} \sum_{s=0}^{\infty} \frac{s!}{(k s)!} a_{r} b_{s} z^{r+k s} .
$$

Since $\bar{\phi} W_{k}^{\star} \bar{\psi}=\bar{\psi} W_{k}^{\star} \bar{\phi}$, therefore it follows that

$$
\sum_{r=0}^{\infty} \sum_{s=0}^{\infty} \frac{s!}{(k s)!} a_{r} b_{s} z^{r+k s}=\sum_{s=0}^{\infty} \sum_{r=0}^{\infty} \frac{r!}{(k r)!} b_{s} a_{r} z^{s+k r}
$$

or equivalently,

$$
\begin{equation*}
\sum_{s=0}^{\infty} \sum_{p=k s}^{\infty} \frac{s!}{(k s)!} a_{p-k s} b_{s} z^{p}=\sum_{r=0}^{\infty} \sum_{p=k r}^{\infty} \frac{r!}{(k r)!} a_{r} b_{p-k r} z^{p} \tag{4.1}
\end{equation*}
$$

For $p=0$, equation (4.1) gives that $a_{0} b_{0}=b_{0} a_{0}$, so $b_{0}=\left(b_{0} / a_{0}\right) a_{0}$, since $a_{0} \neq 0$. Now take $\lambda=b_{0} / a_{0}$, so then $b_{0}=\lambda a_{0}$.
For $1 \leq p \leq k-1$, then by equation (4.1) it follows that $a_{p} b_{0}=b_{p} a_{0}$. This means $b_{p}=\left(b_{0} / a_{0}\right) a_{p}$, since $a_{0} \neq 0$, that is $b_{p}=\lambda a_{p}$.
For $k \leq p \leq 2 k-1$, then by equation (4.1) it follows that $a_{p} b_{0}+\frac{1}{(k)!} a_{p-k} b_{1}=$ $b_{p} a_{0}+\frac{1}{k!} b_{p-k} a_{1}$, but $b_{m}=\lambda a_{m}$ for each $m$ such that $0 \leq m \leq k-1$, therefore $b_{p}=\lambda a_{p}$, since $a_{0} \neq 0$. So continuing in this manner it follows that $b_{i}=\lambda a_{i}$ for each non-negative integer $i$, where $\lambda=b_{0} / a_{0}$. Therefore in this case we get $\bar{\psi}(z)=\sum_{s=0}^{\infty} b_{s} z^{s}=\sum_{s=0}^{\infty} \lambda a_{s} z^{s}=\lambda \bar{\phi}(z)$.
Case III. If $\phi$ and $\psi$ are both not zero identically and $\bar{\phi}(0)=0$ or $\bar{\psi}(0)=0$. Without loss of generality, suppose that $\bar{\phi}^{(m)}(0)=0$ for any integer $m$ such that $0 \leq m \leq m_{1}$ and $\bar{\phi}^{(m+1)}(0) \neq 0$, where $m_{1}$ is a non-negative integer. Therefore by Lemma 4.1, it follows that $\bar{\psi}^{(m)}(0)=0$ for any integer $m$ such that $0 \leq m \leq m_{1}$ and $\bar{\psi}^{(m+1)}(0) \neq 0$. Now we can write $\bar{\phi}, \bar{\psi}$ as

$$
\begin{equation*}
\bar{\phi}(z)=z^{m_{1}+1} \phi_{1}(z), \quad \bar{\psi}(z)=z^{m_{1}+1} \psi_{1}(z) \tag{4.2}
\end{equation*}
$$

where $\phi_{1}, \psi_{1}$ are analytic functions in $L^{\infty}(\mathbb{C})$ and $\phi_{1}(0) \neq 0, \psi_{1}(0) \neq 0$. Therefore by equation (4.2) it follows that,

$$
\bar{\psi} W_{k}^{\star}(\bar{\phi}(z))=z^{(k+1)\left(m_{1}+1\right)} \psi_{1} W_{k}^{\star}\left(\phi_{1}(z)\right)
$$

and

$$
\bar{\phi} W_{k}^{\star}(\bar{\psi}(z))=z^{(k+1)\left(m_{1}+1\right)} \phi_{1} W_{k}^{\star}\left(\psi_{1}(z)\right)
$$

Now as $\bar{\psi} W_{k}^{\star}(\bar{\phi})=\bar{\phi} W_{k}^{\star}(\bar{\psi})$, so it gives $\psi_{1} W_{k}^{\star}\left(\phi_{1}\right)=\phi_{1} W_{k}^{\star}\left(\psi_{1}\right)$. Since $\phi_{1}(0) \neq$ 0 and $\psi_{1}(0) \neq 0$, so by the second case it follows that $\psi_{1}=\lambda_{1} \phi_{1}$, where $\lambda_{1}=\psi_{1}(0) / \phi_{1}(0)$. Therefore we have

$$
\bar{\psi}(z)=z^{m_{1}+1} \psi_{1}(z)=\lambda_{1} z^{m_{1}+1} \phi_{1}(z)=\lambda_{1} \bar{\phi}(z) .
$$

Thus the symbol functions $\phi$ and $\psi$ are linearly dependent.
In ([12]), C. Liu and Yufeng Lu discussed the commutativity of $k$ th-order slant Toeplitz operators on Bergman space and gave the necessary and sufficient conditions for commutativity of $k$ th-order slant Toeplitz operators. In the next theorem, we show that the two $k$ th-order slant Toeplitz operators having harmonic symbols commute if and only if their symbol functions are linearly dependent.

Theorem 4.3. Let $\phi(z)=\sum_{i=0}^{n} a_{i} z^{i}+\sum_{i=1}^{n} a_{-i} \bar{z}^{i}$ and $\psi(z)=\sum_{j=0}^{n} b_{j} z^{j}+$ $\sum_{j=1}^{n} b_{-j} \bar{z}^{j}$, where $b_{-n} \neq 0$ such that the ratio $\frac{\overline{a-n}}{\overline{b_{-n}}}$ is real and $n \geq 1$ be an integer, then the following statements are equivalent:
(1) $B_{\phi}^{k}$ and $B_{\psi}^{k}$ commute;
(2) there exist scalers $\alpha$ and $\beta$, not both zero, such that $\alpha \phi+\beta \psi=0$.

Proof. Suppose that (2) holds then $B_{\phi}^{k}$ and $B_{\psi}^{k}$ commute.
Now suppose that (1) holds. Let $\phi_{1}(z)=\sum_{i=0}^{n} a_{i} z^{i}, \overline{\phi_{2}}(z)=\sum_{i=1}^{n} a_{-i} \bar{z}^{i}$, $\psi_{1}(z)=\sum_{j=0}^{n} b_{j} z^{j}$ and $\overline{\psi_{2}}(z)=\sum_{j=1}^{n} b_{-j} \bar{z}^{j}$. Then $\phi=\phi_{1}+\overline{\phi_{2}}$ and $\psi=\psi_{1}+\overline{\psi_{2}}$. Since $B_{\phi}^{k}$ and $B_{\psi}^{k}$ commute, therefore it gives $T_{\bar{\phi}} W_{k}^{\star} T_{\bar{\psi}} W_{k}^{\star}(1)=T_{\bar{\psi}} W_{k}^{\star} T_{\bar{\phi}} W_{k}^{\star}(1)$, that is

$$
\begin{equation*}
\phi_{2} W_{k}^{\star} \psi_{2}+\overline{\psi_{1}}(0) \phi_{2}+P\left(\overline{\phi_{1}} W_{k}^{\star} \psi_{2}\right)=\psi_{2} W_{k}^{\star} \phi_{2}+\overline{\phi_{1}}(0) \psi_{2}+P\left(\overline{\psi_{1}} W_{k}^{\star} \phi_{2}\right) \tag{4.3}
\end{equation*}
$$

or, equivalently,

$$
\begin{aligned}
& \left(\sum_{i=1}^{n} \overline{a_{-i}} z^{i}\right)\left(\sum_{j=1}^{n} \overline{b_{-j}} \frac{j!}{(k j)!} z^{k j}\right)+\overline{b_{0}} \sum_{i=1}^{n} \overline{a_{-i}} z^{i}+P\left(\left(\sum_{i=0}^{n} \overline{a_{i}} \bar{z}^{i}\right)\left(\sum_{j=1}^{n} \overline{b_{-j}} \frac{j!}{(k j)!} z^{k j}\right)\right) \\
= & \left(\sum_{j=1}^{n} \overline{b_{-j}} z^{j}\right)\left(\sum_{i=1}^{n} \overline{a_{-i}} \frac{i!}{(k i)!} z^{k i}\right)+\overline{a_{0}} \sum_{j=1}^{n} \overline{b_{-j}} z^{j}+P\left(\left(\sum_{j=0}^{n} \overline{b_{j}} \bar{z}^{j}\right)\left(\sum_{i=1}^{n} \overline{a_{-i}} \frac{i!}{(k i)!} z^{k i}\right)\right) .
\end{aligned}
$$

Now for any integer $r$ such that $k n+1 \leq r \leq k n+n$, it follows that

$$
\begin{equation*}
\sum_{i=1}^{n} \sum_{\substack{j=1 \\ i+k j=r}}^{n} \overline{a_{-i}} \overline{b_{-j}} \frac{j!}{(k j)!} z^{r}=\sum_{j=1}^{n} \sum_{\substack{i=1 \\ j+k i=r}}^{n} \overline{b_{-j}} \overline{a_{-i}} \frac{i!}{(k i)!} z^{r} \tag{4.4}
\end{equation*}
$$

where $i$ and $j$ are positive integers not greater than $n$. Now we apply induction to complete the proof.
When $r=k n+n$, then by equation (4.4) it follows that

$$
\overline{a_{-n}} \overline{b_{-n}} \frac{n!}{(k n)!}=\overline{b_{-n}} \overline{\overline{a_{-n}}} \frac{n!}{(k n)!},
$$

so $\overline{a_{-n}}=\overline{b_{-n}}\left(\overline{a_{-n}} / \overline{b_{-n}}\right)$, since $\overline{b_{-n}} \neq 0$. Let $\gamma=\overline{a_{-n}} / \overline{b_{-n}}$, so we get $\overline{a_{-n}}=\gamma \overline{b_{-n}}$. When $r=k n+n-1$, then by equation (4.4) it follows that

$$
\overline{a_{-n+1}} \overline{b_{-n}} \frac{n!}{(k n)!}=\overline{b_{-n+1}} \overline{a_{-n}} \frac{n!}{(k n)!},
$$

so $\overline{a_{-n+1}}=\gamma \overline{b_{-n+1}}$, where $\gamma=\overline{a_{-n}} / \overline{b_{-n}}$, since $\overline{b_{-n}} \neq 0$. Suppose that $\overline{a_{-n+s}}=$ $\gamma \overline{b_{-n+s}}$ for any integer $s$ with $0 \leq s \leq l<n-1$. Now we consider the connection between $a_{-n+l+1}$ and $b_{-n+l+1}$. When $r=k n+n-l-1$, then by equation (4.4) we get that

$$
\begin{aligned}
& \overline{a_{-n+l+1}} \overline{b_{-n}} \frac{n!}{(k n)!}+\cdots+\overline{a_{-n+l+1-k \lambda}} \overline{b_{-n}} \frac{(n+\lambda)!}{(k(n+\lambda))!} \\
& =\overline{b_{-n+l+1}} \overline{a_{-n}} \frac{n!}{(k n)!}+\cdots+\overline{b_{-n+l+1-k \lambda}} \overline{a_{-n+\lambda}} \frac{(n+\lambda)!}{(k(n+\lambda))!}
\end{aligned}
$$

where $\lambda=\left[\frac{(l+1)}{k}\right]$ and $[x]$ denotes the greatest integer function not greater than $x$. Now from the assumption it follows that

$$
\overline{a_{-n+l+1}} \overline{b_{-n}} \frac{n!}{(k n)!}=\overline{b_{-n+l+1}} \overline{a_{-n}} \frac{n!}{(k n)!},
$$

which implies that $\overline{a_{-n+l+1}}=\gamma \overline{b_{-n+l+1}}$, where $\gamma=\overline{a_{-n}} / \overline{b_{-n}}$, since $\overline{b_{-n}} \neq 0$.
Thus, from the induction we obtain that $\overline{a_{-n+s}}=\gamma \overline{b_{-n+s}}$ for any integer $s$ such that $0 \leq s \leq n-1$. Hence $\phi_{2}(z)=\sum_{r=1}^{n} \overline{a_{-r}} z^{r}=\sum_{r=1}^{n} \gamma \overline{b_{-r}} z^{r}=\gamma \psi_{2}(z)$. Now since $\phi_{2}(z)=\gamma \psi_{2}(z)$, so by the equation (4.3) it follows that

$$
\begin{equation*}
\overline{\psi_{1}}(0) \phi_{2}+P\left(\overline{\phi_{1}} W_{k}^{\star} \psi_{2}\right)=\overline{\phi_{1}}(0) \psi_{2}+P\left(\overline{\psi_{1}} W_{k}^{\star} \phi_{2}\right) \tag{4.5}
\end{equation*}
$$

Then

$$
\left\langle\overline{\psi_{1}}(0) \phi_{2}+P\left(\overline{\phi_{1}} W_{k}^{\star} \psi_{2}\right), z^{k n}\right\rangle=\left\langle\overline{\phi_{1}}(0) \psi_{2}+P\left(\overline{\psi_{1}} W_{k}^{\star} \phi_{2}\right), z^{k n}\right\rangle
$$

or equivalently,

$$
\left\langle\overline{\phi_{1}} W_{k}^{\star} \psi_{2}, z^{k n}\right\rangle=\left\langle\overline{\psi_{1}} W_{k}^{\star} \phi_{2}, z^{k n}\right\rangle,
$$

which further implies that $\overline{b_{-n}} \overline{a_{0}} \pi(k n)!=\overline{a_{-n}} \overline{b_{0}} \pi(k n)$ ! and hence $\overline{a_{0}}=\gamma \overline{b_{0}}$, where $\gamma=\overline{a_{-n}} / \overline{b_{-n}}$, since $\overline{b_{-n}} \neq 0$. Now as $a_{0}=\gamma b_{0}$, so from (4.5) it follows that

$$
\overline{\psi_{1}}(0) \phi_{2}=\overline{\phi_{1}}(0) \psi_{2} \text { and } P\left(\overline{\phi_{1}} W_{k}^{\star} \psi_{2}\right)=P\left(\overline{\psi_{1}} W_{k}^{\star} \phi_{2}\right)
$$

that is, $P\left(\overline{\left(\gamma \psi_{1}-\phi_{1}\right)} \cdot W_{k}^{\star} \psi_{2}\right)=0$. Hence for any integer $u$ such that $k n-n \leq$ $u \leq k n-1$, we have

$$
\left\langle z^{u}, P\left(\overline{\left(\gamma \psi_{1}-\phi_{1}\right)} \cdot W_{k}^{\star} \psi_{2}\right)\right\rangle=0
$$

or equivalently,

$$
\left\langle\left(\gamma \psi_{1}-\phi_{1}\right) z^{u}, W_{k}^{\star} \psi_{2}\right\rangle=0
$$

which on putting the values of $\phi_{1}, \psi_{1}$ and $\psi_{2}$ gives that

$$
\left\langle\sum_{r=0}^{n}\left(\gamma b_{r}-a_{r}\right) z^{r+u}, \sum_{r=1}^{n} \frac{r!}{(k r)!} \overline{b_{-r}} z^{k r}\right\rangle=0 .
$$

Since $a_{0}=\gamma b_{0}$, so we have

$$
\begin{equation*}
\left\langle\sum_{r=1}^{n}\left(\gamma b_{r}-a_{r}\right) z^{r+u}, \sum_{r=1}^{n} \frac{r!}{(k r)!} \overline{b_{-r}} z^{k r}\right\rangle=0 . \tag{4.6}
\end{equation*}
$$

When $u=k n-1$, then equation (4.6) gives that

$$
\left(\gamma b_{1}-a_{1}\right) b_{-n} \pi(k n)!=0
$$

so this gives $a_{1}=\gamma b_{1}$, since $b_{-n} \neq 0$. So, then it follows

$$
\left\langle\sum_{r=2}^{n}\left(\gamma b_{r}-a_{r}\right) z^{r+u}, \sum_{r=1}^{n} \frac{r!}{(k r)!} \overline{b_{-r}} z^{k r}\right\rangle=0 .
$$

Now suppose that $a_{j}=\gamma b_{j}$ for any integer $j$ such that $0 \leq j \leq s$, where $0 \leq s \leq$ $n-1$.Then by equation (4.6), it follows that

$$
\begin{equation*}
\left\langle\sum_{r=s+1}^{n}\left(\gamma b_{r}-a_{r}\right) z^{r+u}, \sum_{r=1}^{n} \frac{r!}{(k r)!} \overline{b_{-r}} z^{k r}\right\rangle=0 \tag{4.7}
\end{equation*}
$$

Now consider the connection between $a_{s+1}$ and $b_{s+1}$. When $u=k n-s-1$, then by equation (4.7), it follows that

$$
\left(\gamma b_{s+1}-a_{s+1}\right) b_{-n} \pi(k n)!=0,
$$

which gives $a_{s+1}=\gamma b_{s+1}$, since $b_{-n} \neq 0$. Thus from the induction we obtain that $a_{j}=\gamma b_{j}$ for any integers $j$ with $0 \leq j \leq n$. Therefore we get $\phi_{1}(z)=$ $\sum_{r=1}^{n} a_{r} z^{r}=\sum_{r=1}^{n} \gamma b_{r} z^{r}=\gamma \psi_{1}(z)$. Also from above we have $\phi_{2}=\gamma \psi_{2}$, so it follows that $\phi=\phi_{1}+\overline{\phi_{2}}=\gamma \psi_{1}+\gamma \overline{\psi_{2}}=\gamma \psi$. Thus the symbol functions $\phi$ and $\psi$ are linearly dependent.

Acknowledgement: Support of CSIR Research Grant to first author [F.No. 09/045(1405)/2015-EMR-I] for carrying out the research work is fully acknowledged. The authors would like to express their sincere gratitude to referees for their insightful and valuable comments and suggestions.

## References

1. H. B. An and R. Y. Jian, Slant Toeplitz operators on Bergman spaces, Acta Math. Sinica (Chin. Ser.) 47 (2004) no. 1, 103-110.
2. S. C. Arora and R. Batra, On generalized slant Toeplitz operators, Indian J. Math. 45 (2003) no. 2, 121-134.
3. S. C. Arora and R. Batra, Generalized slant Toeplitz operators on $H^{2}$, Math. Nachr. 278 (2005) no. 4, 347-355.
4. A. Brown and P.R. Halmos, Algebraic properties of Toeplitz operators, J. Reine Angew. Math. 213 (1964), 89-102.
5. R. G. Douglas, Banach algebra techniques in operator theory, Academic Press, New YorkLondon, 1972.
6. T. N. T. Goodman, C. A. Micchelli, and J. D. Ward, Spectral radius formulas for subdivision operators, Recent advances in wavelet analysis, 335-360, Wavelet Anal. Appl., 3, Academic Press, Boston, MA, 1994.
7. C. Heil, G. Strang, and V. Strela, Approximation by translates of refinable functions, Numer. Math. 73 (1996), no. 1, 75-94.
8. M. C. Ho, Properties of slant Toeplitz operators, Indiana Univ. Math. J. 45 (1996), no. 3, 843-862.
9. M. C. Ho, Spectra of slant Toeplitz operators with continuous symbols., Michigan Math. J. 44 (1997), no. 1, 157-166.
10. M. C. Ho, Adjoints of slant Toeplitz operators, Integral Equations Operator Theory 29 (1997), no. 3, 301-312.
11. M. C. Ho, Adjoints of slant Toeplitz operators II, Integral Equations Operator Theory 41 (2001), no. 2, 179-188.
12. C. Liu and Y. Lu, Product and commutativity of kth-order slant Toeplitz operators, Abstr. Appl. Anal. 2013, Art. ID 473916, 11 pp.
13. G. Strang and V. Strela, Orthogonal multiwavelets with vanishing moments, Optical Engineering 33 (1994), no. 7, 2104-2107.
14. G. Strang and V.Strela Short wavelets and matrix dilation equationsIEEE Trans. Signal Process 43 (1995), no. 4, 108-115.
15. K. Stroethoff, The Berezin transform and operators on spaces of analytic functions, Banach Center Publ. 38 (1997), 361-380.
16. O. Toeplitz, Zur Theorie der quadratischen und bilinearen Formen von unendlichvielen Veränderlichen, Math. Ann. 70 (1911), no.3, 351-376.
17. L. F. Villemoes, Wavelet analysis of refinement equations, SIAM J. Math. Anal. 25 (1994), no. 5, 1433-1460.
18. K. Zhu, Analysis on Fock spaces, Graduate Texts in Mathematics, 263. Springer, New York, 2012.
${ }^{1}$ Department of Mathematics, University of Delhi, Delhi-110007, India.
E-mail address: shivamkumarsingh14@gmail.com
${ }^{2}$ Department of Mathematics, Delhi College of Arts and Commerce, UniverSity of Delhi, Delhi-110023, India.

E-mail address: dishna2@yahoo.in


[^0]:    Copyright 2016 by the Tusi Mathematical Research Group.
    Date: Received: Mar. 2, 2017; Accepted: May 19, 2017.
    *Corresponding author.
    2010 Mathematics Subject Classification. Primary 47B35; Secondary 46E20.
    Key words and phrases. kth-order slant Toeplitz operator, Fock space, Berezin transform.

