ON THE LIMITING BEHAVIOR OF A HARMONIC OSCILLATOR WITH RANDOM EXTERNAL DISTURBANCE

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(Received October, 1994; Revised March, 1995)

ABSTRACT

This paper deals with the limiting behavior of a harmonic oscillator under the external random disturbance that is a process of the white noise type. Influence of noises is investigated in resonance and non-resonance cases.

Key words: Harmonic Oscillator, Instantaneous Energy, Differential Equation of the Second Order, Itô Stochastic Differential Equation.

AMS (MOS) subject classifications: 60H10.

1. Introduction

We investigate the harmonic oscillator as a system of motion described by a linear differential equation of the second order

 $m\ddot{u}(t) + ku(t) = q(t)$ while m > 0 and k > 0,

where q(t) is an external disturbance force. In the case, where q(t) is a nonrandom periodic function, the instantaneous energy of the oscillator $\varepsilon(t) = \frac{1}{2}[ku^2(t) + m\dot{u}^2(t)]$ is bounded if the period of the function q(t) is not equal to $2\pi\sqrt{m/k}$ and $\varepsilon(t) \sim t^2$ as $t \to \infty$ if period of function q(t) is equal to $2\pi\sqrt{m/k}$ (resonance).

A model of the random harmonic oscillator with $\varepsilon(t) \sim t$ as $t \to \infty$ was considered by Papanicolau [8] for the case when q(t) is a stationary random process; a model in which $\ln \varepsilon(t) \sim t \to \infty$ was considered by Bendersky and Pastur [1] for the case when q(t) = 0 and k = k(t) is a stationary random process; a model in which $\varepsilon(t) \sim \sqrt{t}$ as $t \to \infty$ was considered by Kulinich [7] for the case when $q(t) = g(w(t))\dot{w}(t)$, with $\dot{w}(t)$ as a "white" noise, g(x) a nonrandom function and $g^2(x)$ integrable over \mathbb{R} .

In the present paper, we consider the external random disturbance of the type $q(t) = f(t)g(w(t))\dot{w}(t)$, where f(t) and g(x) are nonrandom functions and $f^2(t)$ is a periodic function with the period 2L.

The limiting behavior (for $t\to\infty$) of the joint distribution of the random variables $(u(t), \dot{u}(t))$ the distribution of the random variable $\varepsilon(t)$ is investigated in the following cases:

1) $2L \neq 2\pi \sqrt{m/k}$; 2) $2L = 2\pi \sqrt{m/k}$.

It is shown in particular that $\varepsilon(t) \sim t$ if $g^2(x) \sim b \neq 0$ as $|x| \to \infty$ (Theorem 1) and Printed in the U.S.A. ©1995 by North Atlantic Science Publishing Company 265

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 $E\varepsilon(t) \sim t^{\frac{\alpha+1}{2}}$ if $g^2(x) \sim b(x) |x|^{\alpha-1}$, while $\alpha > 0$ and $b(x) = b_1$ for x > 0, and $b(x) = b_2$ for x < 0 (corollary of Theorem 2).

Let u(t) be the distance of a particle from its equilibrium position. We assume that the particle has mass m and that it is fastened to an immobile support by a spring with the coefficient of stiffness k. Then u(t) satisfies the following equation:

$$m\ddot{u}(t) + ku(t) = q(t)$$
 while $u(0) = u_0$ and $\dot{u}(0) = \dot{u}_0$ ($\dot{u} \equiv \frac{d}{dt}u$). (1)

Here q(t) is an external force, u_0 is an initial position and \dot{u}_0 is an initial velocity of the particle. We assume, then, that $u_0 = 0$, $\dot{u}_0 = 0$ and $q(t) = f(t)g(w(t))\dot{w}(t)$, where w(t) is a Wiener process. In this case, equation (1) can be considered as a system of stochastic Itô equations:

$$md\dot{u}(t) = -ku(t)dt + f(t)g(w(t))dw(t)$$

$$du(t) = \dot{u}(t)dt$$
(2)

Lemma: Let function f(t) satisfy the condition, $|\int_{0}^{t} f(s)ds| \leq C$, for every finite t, and let g(x) have the second derivative g''(x) almost everywhere while $\int_{0}^{x} |g''(v)| dv = o(|x|^{\alpha})$ as $|x| \to \infty$ with $\alpha > 0$. Then,

$$\lim_{t\to\infty}t^{-\frac{\alpha+1}{2}}E\mid\int_0^t f(s)g(w(s))ds\mid=0,$$

where w(t) is a Wiener process.

Proof: Since the functions f(t) and g''(x) are integrable over every bounded domain, because of Krylov [5], we can apply Itô's formula to the process $\Phi(t, w(t))$, where $\Phi(t, x) = \int_{0}^{t} f(s) dsg(x)$, and obtain

$$\int_{0}^{t} f(s)g(w(s))ds = \int_{0}^{t} f(s)dsg(w(t)) - \int_{0}^{t} \left[\int_{0}^{s} f(s_{1})ds_{1}\right]g'(w(s))dw(s)$$
$$-\frac{1}{2}\int_{0}^{t} \left[\int_{0}^{s} f(s_{1})ds_{1}\right]g''(w(s))ds = I_{1}(t) + I_{2}(t) + I_{3}(t).$$

It is easy to see that the following inequalities hold true:

$$t^{-\frac{\alpha+1}{2}}E | I_{1}(t) | \leq Ct^{-\frac{\alpha+1}{2}}E | g(w(t)) |$$

$$E(t^{-\frac{\alpha+1}{2}}I_{2}(t))^{2} = t^{-(\alpha+1)}E \int_{0}^{t} [\int_{0}^{s} f(s_{1})ds_{1}g'(w(s))]^{2}ds$$

$$\leq C^{2}t^{-(\alpha+1)}E \int_{0}^{t} [g'(w(s))]^{2}ds$$

$$t^{-\frac{\alpha+1}{2}}E | I_{3}(t) | \leq \frac{1}{2}Ct^{-\frac{\alpha+1}{2}}E \int_{0}^{t} | g''(w(s)) | ds.$$
(3)

Next, applying the Itô formula to the processes $\Phi(w(t))$ and $\Phi_1(w(t))$, where

$$\Phi(x) = 2 \int_{0}^{x} \left[\int_{0}^{z} (g'(v))^{2} dv \right] dz \text{ and } \Phi_{1}(x) = 2 \int_{0}^{x} \left[\int_{0}^{z} |g''(v)| dv \right] dz,$$

we obtain the equations

$$t^{-(\alpha+1)}E \int_{0}^{t} [g'(w(s))]^{2} ds = t^{-(\alpha+1)}E\Phi(w(t))$$
(4)

and

$$t^{-\frac{\alpha+1}{2}}E\int_{0}^{t}|g''(w(s))|\,ds=t^{-\frac{\alpha+1}{2}}E\Phi_{1}(w(t)).$$

The conditions of the Lemma require that $g(x) = o(|x|^{\alpha+1})$, $\Phi(x) = o(|x|^{2\alpha+1})$ and $\Phi_1(x) = o(|x|^{\alpha+1})$. When we take into account that $w(t)t^{-\frac{1}{2}}$ for every t > 0 is standard normal it is easy to ensure that $E\frac{|g(w(t))|}{t^{\frac{\alpha+1}{2}}} \to 0$, $E\frac{\Phi(w(t))}{t^{(\alpha+1)}} \to 0$ and $E\frac{\Phi_1(w(t))}{t^{\frac{\alpha+1}{2}}} \to 0$ as $t \to \infty$. These con-

vergences along with (3) and (4) yield the Lemma.

In what follows, we assume that f(t) in the equations is a continuously differentiable function and that $f^{2}(t)$ has period 2L. Let us denote

$$\begin{aligned} a_0 &= \frac{1}{4L} \int_0^{2L} f^2(t) dt, \quad c_0 = \frac{1}{4L} \int_0^{2L} f^2(t) \cos\left(2\sqrt{k/m} \ t\right) dt, \\ a_1 &= a_0 + c_0, \quad a_2 = a_0 - c_0 \quad \text{and} \quad a_3 = \frac{1}{4L} \int_0^{2L} f^2(t) \sin\left(2\sqrt{k/m} \ t\right) dt. \end{aligned}$$

Theorem 1: Let the function g(x) in equation (2) have a second derivative with

$$\lim_{\|x\| \to \infty} \frac{1}{x} \int_{0}^{x} g^{2}(v) dv = b \quad and \quad \lim_{\|x\| \to \infty} \frac{1}{x} \int_{0}^{x} |g'(v) + g(v)g''(v)| dv = 0.$$

1. Suppose $2L \neq n\pi \sqrt{m/k}$ for any n = 1, 2, ... or $2L = n_0 \pi \sqrt{m/k}$ and at the same time, $c_0 = 0$ and $a_3 = 0$. Then the following hold:

a) The joint distribution of the random variables $(u(t)/\sqrt{t}, \dot{u}(t)/\sqrt{t})$, as $t \to \infty$, converges to the distribution of $(\sqrt{\frac{a_0b}{km}}\zeta_1, \frac{\sqrt{a_0b}}{m}\zeta_2)$, where ζ_1 and ζ_2 are independent standard normal random variables.

b) The distribution of the random variable $t^{-1}\varepsilon(t)$, as $t\to\infty$, converges to the exponential distribution with the parameter $m(a_0b)^{-1}$.

2. Suppose $2L = n_0 \pi \sqrt{m/k}$ and that $c_0 \neq 0$ or $a_3 \neq 0$. Then the following hold:

a)
$$P\{\frac{u(t)}{\sqrt{t}} < x_1, \frac{u(t)}{\sqrt{t}} < x_2\} - F_t(x_1, x_2) \rightarrow 0$$
, where for each $t > 0$, $F_t(x_1, x_2)$ is bivariate

normal with the density:

$$\rho_t(x_1, x_2) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-r^2}} \exp\{-\frac{1}{2(1-r^2)}[Ax_1^2 - 2Bx_1x_2 + Cx_2^2]\},\tag{5}$$

where

$$\begin{split} A &= \frac{\sin^2 \alpha}{\sigma_1^2} + 2r \frac{\sin \alpha \cos \alpha}{\sigma_1 \sigma_2} + \frac{\cos^2 \alpha}{\sigma_2^2}, \quad B = -\frac{\sin \alpha \cos \alpha}{\sigma_1^2} + r \frac{\sin^2 \alpha - \cos^2 \alpha}{\sigma_1 \sigma_2} + \frac{\sin \alpha \cos \alpha}{\sigma_2^2}, \\ C &= \frac{\cos^2 \alpha}{\sigma_1^2} - 2r \frac{\sin \alpha \cos \alpha}{\sigma_1 \sigma_2} + \frac{\sin^2 \alpha}{\sigma_2^2}, \quad r = \frac{a_3}{\sqrt{a_1 a_2}}, \ \sigma_1^2 = a_1 b, \ \sigma_2^2 = a_2 b \text{ and} \\ \alpha &= \sqrt{k/m} t. \end{split}$$

b) The distribution of the random variable $t^{-1}\varepsilon(t)$ converges to the distribution with the density:

$$\rho(x) = \frac{m}{b\sqrt{a_1a_2 - a_3^2}} \exp\{-\frac{xm(a_1 + a_2)}{2b(a_1a_2 - a_3^2)}\} \otimes$$

$$I_0\left(\frac{xm}{b(a_1a_2 - a_3^2)}\sqrt{\frac{1}{4}(a_1 - a_2)^2 + a_3^2}\right), \quad x > 0, \quad (6)$$

$$d \text{ Bessel function of the first kind with zero index and } \rho(x) = 0, \text{ when}$$

where $I_0(x)$ is the modified Bessel function of the first kind with zero index and $\rho(x) = 0$, when x < 0.

Proof: We can write the solution of equation (2) in explicit form [2]:

$$u(t) = \frac{1}{\sqrt{km}} \int_0^t f(s)g(w(s))\sin(\sqrt{k/m}(t-s))dw(s)$$
$$\dot{u}(t) = \frac{1}{m} \int_0^t f(s)g(w(s))\cos(\sqrt{k/m}(t-s))dw(s).$$

Let us introduce the parameter $T \geq T_0 > 0$ and denote

$$u_T(t) = u(tT)/\sqrt{T}, \dot{u}_T(t) = \dot{u}(tT)/\sqrt{T}$$
 and $w_T(t) = w(tT)/\sqrt{T}$.

Then,

$$u_T(t) = \frac{1}{\sqrt{km}} \left[\gamma_T^{(1)}(t) \sin(\sqrt{k/m}tT) - \gamma_T^{(2)}(t) \cos(\sqrt{k/m}tT) \right]$$

and

$$\dot{u}_{T}(t) = \frac{1}{m} \Big[\gamma_{T}^{(1)}(t) \cos(\sqrt{k/m} tT) + \gamma_{T}^{(2)}(t) \sin(\sqrt{k/m} tT) \Big], \tag{7}$$

where

$$\gamma_T^{(1)}(t) = \int_0^t g(w_T(s)\sqrt{T})f(sT)\cos(\sqrt{k/m}sT)dw_T(s)$$

 and

$$\gamma_T^{(2)}(t) = \int_0^t g(w_T(s)\sqrt{T})f(sT)\sin(\sqrt{k/m}sT)dw_T(s$$

Since each process $\gamma_T^{(i)}(t)$ for i = 1, 2 is a martingale with respect to the σ -algebra, $\sigma(w_T(s), \sigma(w_T(s)))$

 $s \leq t$), and since each satisfies the Skorohod condition of compactness of random processes [9], we can assume, without loss of generality, that $\gamma_T^{(i)}(t) \rightarrow \gamma^{(i)}(t)$ for i = 1, 2 and $w_T(t) \rightarrow w(t)$ in probability as $T \rightarrow \infty$ at every point t > 0, where w(t) is a Wiener process and each $\gamma^{(i)}(t)$ is a martingale with respect to the σ -algebra $\sigma(w(s), s \leq t)$.

Thus, (7) implies the convergencies

$$u_T(t) - \frac{1}{\sqrt{km}} \left[\gamma^{(1)}(t) \sin(\sqrt{k/m} tT) - \gamma^{(2)}(t) \cos(\sqrt{k/m} tT) \right] \rightarrow 0$$
(8)

and

$$\dot{u}_T(t) - \frac{1}{m} \left[\gamma^{(1)}(t) \cos(\sqrt{k/m} tT) + \gamma^{(2)}(t) \sin(\sqrt{k/m} tT) \right] \rightarrow 0,$$

in probability, as $T \rightarrow \infty$.

Consider now characteristics of martingales:

$$\begin{aligned} \langle \gamma_T^{(1)}(t) \rangle &= \int_0^t g^2(w_T(s)\sqrt{T}) f^2(sT) \cos^2(\sqrt{k/m}sT) ds \\ \langle \gamma_T^{(2)}(t) \rangle &= \int_0^t g^2(w_T(s)\sqrt{T}) f^2(sT) \sin^2(\sqrt{k/m}sT) ds \\ \langle \gamma_T^{(1)}(t), \gamma_T^{(2)}(t) \rangle &= \frac{1}{2} \int_0^t g^2(w_T(s)\sqrt{T}) f^2(sT) \sin(2\sqrt{k/m}sT) ds \end{aligned}$$

Suppose that for the function $f^{2}(t)$ the first assumption of Theorem 1 is satisfied. It is easy to verify that, in this case,

$$f^{2}(t)\cos^{2}(\sqrt{k/m}t) = a_{0} + \alpha_{1}(t), f^{2}(t)\sin^{2}(\sqrt{k/m}t) = a_{0} + \alpha_{2}(t),$$
(9)

and

$$\frac{1}{2}f^2(t)\sin(2\sqrt{k/m}t) = \alpha_3(t),$$

where $a_0 = \frac{1}{4L} \int_0^{2L} f^2(s) ds$, and there is a constant C > 0 that for all $t \ge 0$ satisfies the inequality,

$$|\int_{0}^{t} \alpha_{i}(s)| \leq C, \ i = 1, 2, 3.$$

Then $\langle \gamma_T^{(i)}(t) \rangle = a_0 \int_0^t g^2(w_T(s)\sqrt{T})ds + \int_0^t g^2(w_T(s)\sqrt{t})\alpha_i(sT)ds = I_T(t) + J_T(t).$

Kulinich [6] implies $I_T(t) \rightarrow \beta(t)$ in probability as $T \rightarrow \infty$, where $\beta(t) = a_0 bt$, and due to the Lemma, $E \mid J_T(t) \mid \rightarrow 0$. Therefore, $\langle \gamma_T^{(i)}(t) \rangle \rightarrow a_0 bt$ in probability as $T \rightarrow \infty$ for i = 1, 2. And for the joint characteristic of martingales $\gamma_T^{(1)}$ and $\gamma_T^{(2)}(t)$, we have the equality,

$$\langle \gamma_T^{(1)}(t), \gamma_T^{(2)}(t) \rangle = \int_0^t g^2(w_T(s)\sqrt{T})\alpha_3(sT)ds,$$

which, due to the Lemma, implies the convergence,

$$E \mid \langle \gamma_T^{(1)}(t), \gamma_T^{(2)}(t) \rangle \mid \to 0 \text{ as } t \to \infty.$$

Hence, for characteristics of the limit martingales we have

$$\langle \gamma^{(i)}(t) \rangle = a_0 bt, \ i = 1, 2 \text{ and } \langle \gamma^{(1)}(t), \gamma^{(2)}(t) \rangle = 0.$$
 (10)

It is easy to see that martingales $\gamma^{(1)}(t)$ and $\gamma^{(2)}(t)$ are continuous with probability 1. Therefore, due to [3], there are independent Wiener processes $w^{(1)}(t)$ and $w^{(2)}(t)$ such that

$$\gamma^{(1)}(t) = \sqrt{a_0 b} w^{(1)}(t)$$
 and $\gamma^{(2)}(t) = \sqrt{a_0 b} w^{(2)}(t)$.

Thus, taking into consideration convergencies (8), we have

$$P\{u_{T}(1) < x_{1}, \dot{u}_{T}(1) < x_{2}\}$$
$$-P\left\{\sqrt{\frac{a_{0}b}{km}}\zeta_{T}^{(1)} < x_{1}, \frac{\sqrt{a_{0}b}}{km}\zeta_{T}^{(2)} < x_{2}\right\} \rightarrow 0 \text{ as } T \rightarrow \infty,$$
(11)

where

 $\zeta_T^{(1)} = w^{(1)}(1)\sin(\sqrt{k/m}T) - w^{(2)}(1)\cos(\sqrt{k/m}T)$

and

$$\zeta_T^{(2)} = w^{(1)}(1)\cos(\sqrt{k/m}T) + w^{(2)}(1)\sin\sqrt{k/m}T).$$

Independence of the normally distributed random variables $w^{(1)}(1)$ and $w^{(2)}(1)$ implies that they have a bivariate normal distribution. Hence, due to [4], $\zeta_T^{(1)}$ and $\zeta_T^{(2)}$ are also bivariate normal for every T.

It is easy to verify, that for every T,

$$E\zeta_T^{(i)} = 0, \ D\zeta_T^{(i)} = 1 \text{ and } E\zeta_T^{(1)}\zeta_T^{(2)} = 0.$$

Therefore, the random variables, $\zeta_T^{(1)}$ and $\zeta_T^{(2)}$, are independent standard normal. Convergence (11) yields the proof of statement 1a) of Theorem 1.

Since for instantaneous energy $\varepsilon(t)$ in system (2) we have the equality,

$$T^{-1}\varepsilon(T) = \frac{1}{2m} ([\gamma_T^{(1)}(1)]^2 + [\gamma_T^{(2)}(1)]^2),$$
(12)

then, for all x > 0,

$$\lim_{T \to \infty} P\{T^{-1}\varepsilon(T) < x\} = P\{\frac{a_0 b}{2m}([w^{(1)}(1)]^2 + [w^{(2)}(1)]^2) < x\}$$

According to Gnedenko [4], the random variable $[w^{(1)}(1)]^2 + [w^{(2)}(1)]^2$ has a χ^2 distribution with two degrees of freedom and it coincides with the exponential distribution with parameter 1/2. Hence, the distribution of the random variable $T^{-1}\varepsilon(T)$, as $T\to\infty$ converges to the exponential distribution with parameter $m[a_0b]^{-1}$. This proves statement 1b) of Theorem 1.

Next, suppose that $2L = n_0 \pi \sqrt{m/k}$ and that at least one of the constants, c_0 or a_3 , is not equal to zero. Then (9) can be represented in the form:

$$f^{2}(t) \cos^{2}(\sqrt{k/m}t) = a_{1} + \alpha_{1}(t), f^{2}(t) \sin^{2}(\sqrt{k/m}t) = a_{2} + \alpha_{2}(t)$$
$$\frac{1}{2}f^{2}(t)\sin(2\sqrt{k/m}t) = a_{3} + \alpha_{3}(t).$$

and

Therefore, in this case we have

$$\langle \gamma_T^{(i)}(t) \rangle = a_i \int_0^t g^2(w_T(s)\sqrt{T})ds + \int_0^t g^2(w_T(s)\sqrt{T})\alpha_i(sT)ds, \quad i = 1, 2,$$

and

$$\langle \gamma_T^{(1)}(t), \gamma_T^{(2)}(t) \rangle = a_3 \int_0^t g^2(w_T(s)\sqrt{T})ds + \int_0^t g^2(w_T(s)\sqrt{T})\alpha_3(sT)ds.$$

As in the proof of statement 1 of Theorem 1, we obtain characteristics of the limit martingales:

$$\langle \gamma^{(1)}(t) \rangle = a_1 bt, \ \langle \gamma^{(2)}(t) \rangle = a_2 bt \text{ and } \langle \gamma^{(1)}(t), \gamma^{(2)}(t) \rangle = a_3 bt.$$

Also,

$$\gamma^{(i)}(t) = \sqrt{b}[b_{i1}w^{(1)}(t) + b_{i2}w^{(2)}(t)], \quad i = 1, 2,$$

where $w^{(1)}(t)$ and $w^{(2)}(t)$ are independent Wiener processes and (b_{i1}, b_{i2}) is the *i*-th row of the matrix $B^{1/2}$, where $B = \begin{pmatrix} a_1 & a_3 \\ a_3 & a_2 \end{pmatrix}$. The independence of the Wiener processes, $w^{(1)}(t)$ and $w^{(2)}(t)$, implies that random variables $\gamma^{(1)}(1)$ and $\gamma^{(2)}(1)$ have normal distributions with parameters $(0, \sigma_i^2)$, where $\sigma_1^2 = a_1 b$ and $\sigma_2^2 = a_2 b$ are bivariate normal with the coefficient of correlation $r = a_3(a_1a_2)^{-1}$. Hence, according to Gnedenko [4], the joint density of the random variables,

$$\gamma^{(1)}(1) \sin(\sqrt{k/m}T) - \gamma^{(2)}(1) \cos(\sqrt{k/m}T)$$

and

$$\gamma^{(1)}(1)\cos(\sqrt{k/m}T) + \gamma^{(2)}(1)\sin\sqrt{k/m}T),$$

is of the form (5) with t = T. To complete the proof of statement 2*a*) of Theorem 1, we use convergencies (8). Equality (12) implies that the limit distribution of the random variable $T^{-1}\varepsilon(T)$ coincides with the distribution of the absolute value of a bivariate normal random vector.

Corollary: Under the conditions of Theorem 1,

$$\begin{split} \lim_{t \to \infty} (Et^{-1}\varepsilon(t)) &= \frac{b}{2m}(a_1 + a_2);\\ \lim_{t \to \infty} Dt^{-1}\varepsilon(t) &= \frac{b^2}{2m^2}(a_1^2 + a_2^2), \text{ while } a_3 = 0;\\ \lim_{t \to \infty} Dt^{-1}\varepsilon(t) &= \frac{b^2}{2m^2}(a_0^2 + \frac{3}{2}a_3^2), \text{ while } a_3 \neq 0 \text{ and } c_0 = 0;\\ \lim_{t \to \infty} Dt^{-1}\varepsilon(t) &= \frac{b^2}{2m^2}(a_1^2 + a_2^2 + 3a_3^2 + \beta), \text{ while } a_3 \neq 0 \quad c_0 \neq 0 \text{ and}\\ \beta &= \frac{2a_3^2}{(a_1 - a_2)^4} \{4a_3^2[a_1^4 - a_1^2\sqrt{a_1^4 - (a_1 - a_2)^2} - \frac{1}{2}(a_1 - a_2)^2]\\ &- [a_1^4 - (a_1^2 - a_2^2)\sqrt{a_1^4 - (a_1 - a_2)^2}]\} + 2(a_1a_2a_3)^2 + a_3^2. \end{split}$$

In this case we can change the order of limit and expectation (variance). We use the latter, the explicit form of the limit value $\gamma^{(i)}(1)$ for every *i* and equality (12) to prove the statement.

Theorem 2: Let the function g(x) in equation (2) have a second derivative almost everywhere and for some $\alpha > 0$ satisfy the conditions:

$$\lim_{\|x\| \to \infty} \left(\frac{1}{\|x\|^{\alpha}} \int_{0}^{x} g^{2}(v) dv - b(x) \right) = 0, \text{ with } b(x) = \begin{cases} b_{1}, & x > 0 \\ b_{2}, & x < 0 \end{cases}$$

and

$$\lim_{\|x\|\to\infty} \frac{1}{\|x\|^{\alpha}} \int_{0}^{x} |g'(v) + g(v)g''(v)| dv = 0.$$

Then

$$P\left\{\frac{u(t)}{t^{(\alpha+1)/4}} < x_1, \frac{\dot{u}(t)}{t^{(\alpha+1)/4}} < x_2\right\} - P\{v(t) < x_1, \dot{v}(t) < x_2\} \rightarrow 0 \text{ as } t \rightarrow \infty,$$

where v(t) is the position and $\dot{v}(t)$ is the velocity of the homogeneous harmonic oscillator

$$m\ddot{v}(t) + kv(t) = 0, \quad t > 0$$
 (13)

with the initial condition

$$v(0) = -\frac{1}{\sqrt{km}}\gamma^{(2)}(1)$$
 and $\dot{v}(0) = \frac{1}{m}\gamma^{(1)}(1)$.

Here each $\gamma^{(i)}(t)$ is a martingale with respect to the σ -algebra $\sigma(w(s), s \leq t)$ with characteristics:

$$\langle \gamma^{(i)}(t) \rangle = a_i \beta(t), \quad i = 1, 2 \text{ and } \langle \gamma^{(1)}(t), \gamma^{(2)}(t) \rangle = a_3 \beta(t),$$

while $\beta(t) = \alpha \int_0^t |w(s)|^{\alpha - 1} b(w(s)) \operatorname{sign} w(s) ds.$

Proof: The proof is similar to that of (8) in Theorem 1, with the difference that, in this case,

$$u_T(t) = T^{-(\alpha+1)/4} u(tT), \quad \dot{u}_T(t) = T^{-(\alpha+1)/4} \dot{u}(tT),$$

$$\gamma_T^{(1)}(t) = T^{(1-\alpha)/4} \int_0^t g(w_T(s)\sqrt{T}) f(sT) \cos(\sqrt{k/m}sT) dw_T(s),$$

and

$$\gamma_T^{(2)}(t) = T^{(1-\alpha)/4} \int_0^t g(w_T(s)\sqrt{T})f(sT)\sin(\sqrt{k/m}sT)dw_T(s)$$

with characteristics

$$\langle \gamma_T^{(i)}(t) \rangle = a_i T^{(1-\alpha)/2} \int_0^t g^2(w_T(s)\sqrt{T}) ds + T^{(1-\alpha)/2} \int_0^t g^2(w_T(s)\sqrt{T})\alpha_i(sT) ds, \ i = 1, 2,$$

and

$$\langle \gamma_T^{(1)}(t), \gamma_T^{(2)}(t) \rangle = a_3 T^{(1-\alpha)/2} \int_0^t g^2(w_T(s)\sqrt{T}) ds + T^{(1-\alpha)/2} \int_0^t g^2(w_T(s)\sqrt{T})\alpha_3(sT) ds.$$

Due to the Lemma,

$$\langle \gamma_T^{(i)}(t) \rangle = a_i T^{(1-\alpha)/2} \int_0^t g^2(w_T(s)\sqrt{T}) ds + o(1)$$

and

$$\langle \gamma_T^{(1)}(t), \gamma_T^{(2)}(t) \rangle = a_3 T^{(1-\alpha)/2} \int_0^t g^2(w_T(s)\sqrt{T}) ds + o(1),$$

where o(1) is such that $E \mid o(1) \mid \to 0$ as $T \to \infty$ for all t > 0. Next, Kulinich [6] established that

$$T^{(1-\alpha)/2} \int_{0}^{t} g^2(w_T(s)\sqrt{T}) ds \rightarrow \beta(t)$$

in probability as $T \rightarrow \infty$, where

$$\beta(t) = 2\left[\int_{0}^{w(t)} |v|^{\alpha} b(v) dv - \int_{0}^{t} |w(s)|^{\alpha} b(w(s)) dw(s)\right].$$

Since $\alpha > 0$, using Itô's formula, we have

$$\beta(t) = \alpha \int_0^t |w(s)|^{\alpha - 1} b(w(s)) \operatorname{sign} w(s) ds.$$
(14)

Hence,

$$\langle \gamma_T^{(i)}(t) \rangle \rightarrow a_i \beta(t), \ i = 1, 2 \text{ and } \langle \gamma_T^{(1)}(t), \gamma_T^{(2)}(t) \rangle \rightarrow a_3 \beta(t)$$

in probability as $T \rightarrow \infty$. Thus, we obtain convergence (8), where each $\gamma^{(i)}(t)$ is a continuous, with probability 1, martingale with respect to $\sigma(w(s), s \leq t)$, with characteristics:

$$\langle \gamma^{(i)}(t) \rangle = a_i \beta(t), i = 1, 2 \text{ and } \langle \gamma^{(1)}(t)(t), \gamma^{(2)}(t) \rangle = a_3 \beta(t),$$

where $\beta(t)$ has the form (14). Using convergence (8) for t = 1 and an explicit form of the solution of problem (13), we complete the proof of Theorem 2.

Corollary: Under the conditions of Theorem 2,

$$\lim_{t \to \infty} Et^{-(\alpha+1)/2} \varepsilon(t) = \frac{\alpha(a_1+a_2)}{2m} \int_0^1 E |w(s)|^{\alpha-1} b(w(s)) \operatorname{sign} w(s) ds.$$

This equality is a consequence of the following statements:

- 1) the equality (12); 2) the equality $E[\gamma_T^{(i)}(t)]^2 = E\langle \gamma_T^{(i)}(t) \rangle$ 3) the possibility to change the order of limit and expectation.

Remark: Let $q(x_1, x_2)$ be a joint density of the distribution of $\gamma^{(1)}(1)$ and $\gamma^{(2)}(1)$ and $\rho_t(x_1, x_2)$ be a joint density of the distribution of the position v(t) and the velocity $\dot{v}(t)$ at the moment t, described by (13). Then,

$$\rho_t(x_1, x_3) = q[x_1\sqrt{km}\sin\left(\sqrt{k/mt}\right) + x_2m\cos\left(\sqrt{k/mt}\right),$$
$$-x_1\sqrt{km}\cos\left(\sqrt{k/mt}\right) + x_2m\sin\left(\sqrt{k/mt}\right)]m\sqrt{km}.$$
(15)

Using the explicit form of the solution to equation (13) we get

$$\gamma^{(1)}(1) = v(t)\sqrt{km}\sin\left(\sqrt{k/m}t\right) + \dot{v}(t)m\cos\left(\sqrt{k/m}t\right)$$

and

$$\gamma^{(2)}(1) = -v(t)\sqrt{km}\cos\left(\sqrt{k/m}t\right) + \dot{v}(t)m\sin\left(\sqrt{k/m}t\right)$$

which yields (15).

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