QUANTITATIVE RESULTS FOR PERTURBED STOCHASTIC DIFFERENTIAL EQUATIONS

JORDAN STOYANOV and DOBRIN BOTEV

Bulgarian Academy of Sciences Institute of Mathematics 1090 Sofia, Box 373, Bulgaria

(Received August, 1995; Revised December, 1995)

ABSTRACT

The paper is devoted to Itô type stochastic differential equations (SDE's) with "small" perturbations. Our goal is to present strong results showing how "close" are the 2m-order moments of the solutions of the perturbed SDE's and the unperturbed SDE.

Key words: Stochastic Differential Equation, "Small" Perturbations, Moments of Order 2m, Approximations of the Solutions.

AMS (MOS) subject classifications: 60H10.

1. Introduction. Statement of the Problem

The object of this study are stochastic differential equations (SDE's) of the following type

$$X_{t} = X_{t_{0}} + \int_{t_{0}}^{t} a(s, X_{s})ds + \int_{t_{0}}^{t} b(s, X_{s})dw_{s}, \quad t \ge t_{0} \ge 0.$$
(1)

Here $w = (w_t, t \ge 0)$ is a standard Wiener process defined on a given probability space $(\Omega, \mathfrak{F}, P)$, a(t,x) and b(t,x), $t \ge t_0$, $x \in \mathbb{R}^1$, are measurable real-valued functions, and X_{t_0} is a random variable (r.v.) independent of w with $E\{X_{t_0}^2\} < \infty$. Finally, $\int_{t_0}^t b(\cdot) dw_s$ is the well-known stochastic integral in Itô sense.

Under general conditions, the SDE (1) has a unique (strong) solution $X = (X_t, t \ge t_0)$, which is a diffusion Markov process with a drift coefficient *a* and a diffusion coefficient b^2 . Let us adopt the following classical conditions: For some constants $K_1 > 0$ and $K_2 > 0$ and all $t \ge t_0$, $x, y \in \mathbb{R}^1$ we have

$$|a(t,x) - a(t,y)| + |b(t,x) - b(t,y)| \le K_1 |x - y| a^2(t,x) + b^2(t,x) \le K_2^2(1 + x^2).$$

$$(2)$$

Notice that standard references in the area of SDE's are the books by Gihman and Skorohod [4], Arnold [1], Liptser and Shiryaev [5] and Gard [3]. We shall systematically use basic facts from these sources without mentioning.

Now, along with (1), we consider another SDE:

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$$X_t^{\epsilon} = X_{t_0}^{\epsilon_0} + \int_{t_0}^t \widetilde{a} (s, X_s^{\epsilon}, \epsilon_1) ds + \int_0^t \widetilde{b} (s, X_s^{\epsilon}, \epsilon_2) dw_s, \quad t \ge t_0.$$
(3)

Here ϵ_0 , ϵ_1 , ϵ_2 are "small" positive parameters, e.g., each is in the interval (0,1] and ϵ stands for $(\epsilon_0, \epsilon_1, \epsilon_2)$; $X_{t_0}^{\epsilon_0}$, \tilde{a} , and \tilde{b} are as X_{t_0} , a, and b above; w is the same Wiener process. Thus the SDE (3) has also a unique (strong) solution $X^{\epsilon} = (X_t^{\epsilon}, t \ge t_0)$.

Our goal is to compare the solutions X^{ϵ} and X of (3) and (1) in the case when their coefficients are related as follows:

$$\begin{cases} \widetilde{a}(t, x, \epsilon_1) = a(t, x) + \alpha(t, x, \epsilon_1) \\ \widetilde{b}(t, x, \epsilon_2) = b(t, x) + \beta(t, x, \epsilon_2). \end{cases}$$
(4)

The terms $\alpha(\cdot)$ and $\beta(\cdot)$ are called perturbations of the coefficients $a(\cdot)$ and $b(\cdot)$ which explains why (3) is called a "perturbed SDE" while the name "unperturbed SDE" is kept for (1).

Let us suppose that for some fixed natural number m we have $E\{(X_{t_0}^{\epsilon_0})^{2m}\} < \infty$, $E\{(X_{t_0})^{2m}\} < \infty$ and let for all $t \ge t_0$,

$$E\{ |X_{t_0}^{\epsilon_0} - X_{t_0}|^{2m} \} \leq \delta_0(\epsilon_0)$$

$$\sup_x |\alpha(t, x, \epsilon_1)| \leq \delta_1(t, \epsilon_1)$$

$$\sup_x |\beta(t, x, \epsilon_2)| \leq \delta_2(t, \epsilon_2).$$
(5)

Thus we can expect that if the quantities $\delta_0(\epsilon_0)$, $\delta_1(t,\epsilon_1)$, $\delta_2(t,\epsilon_2)$ are small for small ϵ_0 , ϵ_1 , ϵ_2 , then the process X^{ϵ} is close to X. Recall that X^{ϵ} and X are 2*m*-integrable in the sense that for each $t \geq t_0$ the r.v.'s $|X_t^{\epsilon}|^{2m}$ and $|X_t|^{2m}$ are P-integrable. Thus, the following quantity

$$\Delta_t^{\epsilon} = E\{ \mid X_t^{\epsilon} - X_t \mid {}^{2m} \}$$

is well-defined and we are interested in conditions guaranteeing that $\Delta_t^{\epsilon} \to 0$ as $\epsilon \to 0$. This means that for a fixed t, the 2*m*-order moments of X_t^{ϵ} and X_t are close. Furthermore, we describe a few cases when $\Delta_t^{\epsilon} \to 0$ as $\epsilon \to 0$ on intervals whose length tend to infinity.

2. Preliminary Result

Let us prove first a result which is of independent interest. This result plays a key role for the statements in the next section.

Theorem A: Suppose conditions (2) and (5) are satisfied for the SDE's (1) and (3). Then for any $t \ge t_0$,

$$\Delta_t^{\epsilon} \le \left\{ \delta_0^{1/m}(\epsilon_0) \exp\left[M(t-t_0) + 2\int_{t_0}^t \delta_1(s,\epsilon_1) ds \right] + \int_{t_0}^t [2\delta_1(s,\epsilon_1) + (2m-1)\delta_2^2(s,\epsilon_2)] \exp\left[M(t-s) + 2\int_s^t \delta_1(\tau,\epsilon_1) d\tau \right] ds \right\}^m,$$
(6)

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where $M = 2K + 2(2m-1)K^2$ ($K = \max[K_1, K_2]$, and K_1 and K_2 are the constants from (2)).

Proof: If we write explicitly the difference $Z_t^{\epsilon} = X_t^{\epsilon} - X_t$, $t \ge t_0$, apply the Itô formula to $(Z_t^{\epsilon})^{2m}$ and take expectations, we find that

$$\Delta_t^{\epsilon} = E\{(Z_t^{\epsilon})^{2m}\} = E\{(Z_{t_0}^{\epsilon_0})^{2m}\} + 2mE\{I_1(t)\} + m(2m-1)E\{I_2(t)\} + 2mE\{I_3(t)\},$$
(7)

where

$$\begin{split} I_1(t) &= \int_{t_0}^t [\widetilde{a} \; (s, X_s^{\epsilon}, \epsilon_1) - a(s, X_s)] (Z_s^{\epsilon})^{2m-1} ds, \\ I_2(t) &= \int_{t_0}^t \left[\widetilde{b} \; (s, X_s^{\epsilon}, \epsilon_2) - b(s, X_s) \right]^2 (Z_s^{\epsilon})^{2m-2} ds, \\ I_3(t) &= \int_{t_0}^t \left[\widetilde{b} \; (s, X_s^{\epsilon}, \epsilon_2) - b(s, X_s) \right] (Z_s^{\epsilon})^{2m-1} dw_s. \end{split}$$

The existence of the 2*m*-order moments of X_t^{ϵ} and X_t , $t \ge t_0$ and the conditions on $\tilde{b}(\cdot)$ and $b(\cdot)$ allow us to use one of the properties of the stochastic integrals, thus concluding that

$$E\{I_3(t)\} = 0.$$

Let us estimate $I_1(t)$ and $I_2(t)$. In view of (2), we see that

$$I_1(t) \leq K \int\limits_{t_0}^t \mid X_s^{\epsilon} - X_s \mid \ \mid Z_s^{\epsilon} \mid ^{2m-1}ds + \int\limits_{t_0}^t \delta_1(s,\epsilon_1) \mid Z_s^{\epsilon} \mid ^{2m-1}ds$$

Now we take the expectations of both sides of the last inequality and applying Hölder's inequality to $E\{ |Z_s^{\epsilon}|^{2m-1} \}$ (see Shiryaev [7]) we find that

$$E\{I_1(t)\} \le \int_{t_0}^t \Delta_s^{\epsilon} ds + \int_{t_0}^t \delta_1(s,\epsilon_1) (\Delta_s^{\epsilon})^{(2m-1)/(2m)} ds.$$

$$\tag{8}$$

Similar arguments imply that

$$E\{I_2(t)\} \le 2K^2 \int_{t_0}^t \Delta_s^{\epsilon} ds + \int_{t_0}^t \delta_2^2(s,\epsilon_2) (\Delta_s^{\epsilon})^{(m-1)/m} ds.$$
(9)
and (9) we get

Therefore, from (7), (8) and (9), we get

$$\Delta_t^{\epsilon} \le \delta_0(\epsilon_0) + mM \int_{t_0}^t \Delta_s^{\epsilon} ds + 2m \int_{t_0}^t \delta_1(s,\epsilon_1) (\Delta_s^{\epsilon})^{(2m-1)/(2m)} ds$$
(10)

$$+ m(2m-1) \int_{t_0}^t \delta_2^2(s,\epsilon_2) (\Delta_s^{\epsilon})^{(m-1)/m} ds.$$

Now we use the following elementary inequality

 $v^{r_2} \le v^{r_1} + v$

which is valid for any nonnegative number v and $0 < r_1 \leq r_2 < 1. \ {\rm Setting}$

$$r_1 = (m-1)/m, \quad r_2 = (2m-1)/(2m), \text{ and } v = \Delta_s^\epsilon$$

we find that

$$(\Delta_s^{\epsilon})^{(2m-1)/(2m)} \le (\Delta_s^{\epsilon})^{(m-1)/m} + \Delta_s^{\epsilon}$$

and hence (10) takes the form

$$\Delta_t^{\epsilon} \le \delta_0(\epsilon_0) + m \int_{t_0}^t [M + 2\delta_1(s, \epsilon_1)] \Delta_s^{\epsilon} ds$$

$$+ \int_{t_0}^t [2m\delta_1(s, \epsilon_1) + m(2m - 1)\delta_2^2(s, \epsilon_2)] (\Delta_s^{\epsilon})^{(m-1)/m} ds.$$
(11)

The last tool we need is the following generalized Gronwall-Bellman inequality (see Filatov and Sharova [2]):

If a nonnegative function u(t), $t \ge t_0$, satisfies the integral inequality

$$u(t) \leq C + \int_{t_0}^t A(s)u(s)ds + \int_{t_0}^t B(s)[u(s)]^{\gamma}ds,$$

where $C \ge 0$, $0 \le \gamma < 1$, and functions A(t) and B(t), $t \ge t_0$, are nonnegative and continuous, then

$$u(t) \leq \left\{ C^{1-\gamma} \exp\left[(1-\gamma) \int_{t_0}^t A(s) ds \right] + (1-\gamma) \int_{t_0}^t B(s) \exp\left[(1-\gamma) \int_s^t A(\tau) d\tau \right] ds \right\}^{1/(1-\gamma)}$$

Obviously, it remains to apply this inequality to (11) by letting

$$\begin{split} u(t) &= \Delta_t^{\epsilon}, \quad C = \delta_0(\epsilon_0), \quad \gamma = (m-1)/m, \\ A(s) &= m[M + 2\delta_1(s,\epsilon_1)], \\ B(s) &= 2m\delta_1(s,\epsilon_1) + m(2m-1)\delta_2^2(s,\epsilon_2). \end{split}$$

Thus we arrive at the desired relation (6). Theorem A is proved.

3. Basic Results. Proofs

Since the magnitude of the perturbations of SDE (1) is determined by the quantities $\delta_0(\epsilon_0)$, $\delta_1(t,\epsilon_1)$ and $\delta_2(t,\epsilon_2)$ (see (4) and (5)) it is natural to impose some conditions on these quantities and see how $\Delta_t^{\epsilon} = E\{|X_t^{\epsilon} - X_t|^{2m}\} \rightarrow 0$ as $\epsilon \rightarrow 0$ and on which intervals this convergence holds.

Three specific cases will be considered.

In the statements below (Theorems 1, 2 and 3) we assume (with mentioning it again) that the general conditions of Theorem A are satisfied. We also use the constant $M = 2K + 2(2m-1)K^2$.

Theorem 1: Suppose that for all $t \ge t_0$,

$$\delta_0(\epsilon_0) = \epsilon_0, \ \delta_1(t, \epsilon_1) = \epsilon_1 \ and \ \delta_2(t, \epsilon_2) = \epsilon_2. \tag{12}$$

Define the numbers ϵ and T_1 as follows:

$$\epsilon = \max\left[\epsilon_0^{1/m}, \epsilon_1, \epsilon_2^2\right] \text{ and } T_1 = (1-\rho)/(M+\rho),$$

where $\rho \in (0,1)$ is arbitrarily chosen.

Then the following relation holds:

$$\sup_{t} \Delta_t^{\epsilon} \to 0 \text{ as } \epsilon \to 0 \text{ for } t \in [t_0, t_0 + T_1 \ln(1/\epsilon)).$$

Proof: In Theorem A we have found the upper bound (6) for Δ_t^{ϵ} and now by using (12) we can go further. After a substitution we arrive at

$$(\Delta_t^{\epsilon})^{1/m} \le \epsilon_0^{1/m} \exp\left[M(t-t_0) + 2\epsilon_1(t-t_0)\right] + \int_{t_0}^{t} \left[2\epsilon_1 + (2m-1)\epsilon_2^2 \exp\left[M(t-s) + 2\epsilon_1(t-s)\right]\right] ds.$$

Since $\epsilon \to 0$ implies that $\epsilon_1 \to 0$, we can assume that $2\epsilon_1 < \rho$. Taking into account that $\epsilon_0^{1/m} \leq \epsilon, \ \epsilon_1 \leq \epsilon \text{ and } \epsilon_2^2 \leq \epsilon$ we find that

$$(\Delta_t^{\epsilon})^{1/m} \le C_1 \epsilon e^{(M+\rho)t} + (4\epsilon)/(M+\rho), \tag{13}$$

where the constant C_1 depends on t_0 and m but not on t.

Obviously, (13) implies that $\Delta_t^{\epsilon} \to 0$ as $\epsilon \to 0$ for each t on any finite interval $[t_0, t_1]$ with fixed $t_1 > t_0$. However, we can use (13) and make one step further by extending the time-interval on which $\Delta_t^{\epsilon} \to 0$. Indeed, if we take $T_1 = (1 - \rho)/(M + \rho)$ we find from (13) that

$$\Delta_t^\epsilon \leq [C_1 \epsilon^\rho + (4\epsilon)/(M+\rho)]^m \text{ for any } t \in [t_0,t_0+T_1 \ln{(1/\epsilon)})$$

and hence $\sup_t \Delta_t^{\epsilon} \to 0$ as $\epsilon \to 0$ on the interval $[t_0, t_0 + T_1 \ln(1/\epsilon))$. Note that the length of this interval tends to infinity as $\epsilon \to 0$. Theorem 1 is proved.

Theorem 2: Suppose that $t_0 > 1$ and let for all $t \ge t_0$,

$$\delta_0(\epsilon_0) = t_0^{-1/\epsilon_0}, \ \delta_1(t,\epsilon_1) = t^{-1/\epsilon_1} \ and \ \delta_2(t,\epsilon_2) = t^{-1/\epsilon_2}.$$
(14)

Define ϵ and T_2 as follows:

$$\epsilon = \max\left[m\epsilon_0, \epsilon_1, \epsilon_2/2\right] \text{ and } T_2 = (1/M)\ln(t_0 - \rho),$$

where ρ is an arbitrary number in the interval $(0, t_0)$.

Then,

Proof: From (6) and (14), we find that

$$(\Delta_t^{\epsilon})^{1/m} \le t_0^{-1/(m\epsilon_0)} \exp\left[M(t-t_0) + 2\epsilon_1(t_0^{1-1/\epsilon_1} - t^{1-1/\epsilon_1})/(1-\epsilon_1 0)\right] + \int_{t_0}^t \left[2s^{-1/\epsilon_1} + (2m-1)s^{-2/\epsilon_2}\right] \exp\left[M(t-s) + 2\epsilon_1(s^{1-1/\epsilon_1} - t^{1-1/\epsilon})/(1-\epsilon_1 0)ds.$$

Obviously, we can take $\epsilon_1 < 1/2$, and, since $s^{-1/(m\epsilon_0)} \le s^{-1/\epsilon}$, $s^{-1/\epsilon_1} \le s^{-1/\epsilon}$ and $s^{-2/\epsilon_2} \le s^{-1/\epsilon}$ for any s > 1, we derive that

$$(\Delta_t^{\epsilon})^{1/m} \le t_0^{-1/\epsilon} e^{M(t-t_0)+2} + 4m \int_{t_0}^{t} s^{-1/\epsilon} e^{M(t-s)+2} ds.$$

Hence,

$$(\Delta_t^{\epsilon})^{1/m} \le C_2 t_0^{-1/\epsilon} e^{Mt},\tag{15}$$

where C_2 is a constant depending on t_0 and m but not on t. It follows from (15) that $\Delta_t^{\epsilon} \rightarrow 0$ as $\epsilon \rightarrow 0$ for each t on any finite interval $[t_0, t_2]$ with fixed $t_2 > t_0$. Moreover, with $T_2 = (1/M) \ln(t_0 - \rho)$ we easily find that

$$\Delta_t^{\epsilon} \le C_2^m (1 - \rho/t_0)^{m/\epsilon} \text{ for all } t \in [t_0, t_0 + T_2/\epsilon).$$

Since $0 < 1 - \rho/t_0 < 1$, the conclusion is that $\sup_t \Delta_t^{\epsilon} \to 0$ as $\epsilon \to 0$ on the interval $[t_0, t_0 + T_2/\epsilon)$ whose length tends to infinity. Theorem 2 is proved.

Theorem 3: Suppose that $t_0 > 0$ and let for all $t \ge t_0$

$$\delta_0(\epsilon_0) = e^{-t_0/\epsilon_0}, \ \delta_1(t,\epsilon_1) = e^{-t/\epsilon_1} \ and \ \delta_2(t,\epsilon_2) = e^{-t/\epsilon_2}.$$
(16)

For an arbitrary $\rho \in (0, t_0)$, define

$$\epsilon = \max[m\epsilon_0, \epsilon_2/2]$$
 and $T_3 = (t_0 - \rho)/M_{\odot}$

Then,

$$\sup_t \Delta_t^\epsilon \to 0 \ as \ \epsilon \to 0 \ for \ t \in [t_0, t_0 + T_3/\epsilon).$$

Proof: Taking $\epsilon_1 < 1$ (we can do this since $\epsilon_1 \rightarrow 0$) we easily see that $\int_{t_0}^{t} e^{-s/\epsilon_1} ds < 1$ and $\int_{s}^{t} e^{-\tau/\epsilon_1} d\tau < 1$. Then, substituting (16) into (6) and suitably transforming the right-hand side of (6) we finally arrive at the relation:

$$(\Delta_t^{\epsilon})^{1/m} \le C_3 e^{Mt - t_0/\epsilon},$$

where C_3 depends on t_0 , m and M but not on t.

Therefore, $\Delta_t^{\epsilon} \to 0$ as $\epsilon \to 0$ for each t on any finite interval $[t_0, t_3]$ with fixed $t_3 > t_0$. Even more, if $T_3 = (t_0 = \rho)/M$, then

$$\sup_t \Delta_t^{\epsilon} \leq C_3^m e^{-\rho/\epsilon} \to 0 \text{ as } \epsilon \to 0 \text{ for } t \in [t_0, t_0 + T_3/\epsilon).$$

Again the convergence to zero holds on intervals whose lengths tends to infinity. Theorem 3 is proved. $\hfill \Box$

4. Additional Remarks

(a) In Theorems 1, 2 and 3 not only did we prove that $\sup_t \Delta_t^{\epsilon} \to 0$ as $\epsilon \to 0$ but, in addition, we can specify the rate of convergence. It is a power rate in Theorem 1 and exponential rate in Theorems 2 and 3. Moreover, we can specify the rate of getting to infinity of the lengths of the corresponding intervals. Obviously, both rates depend on the magnitude of perturbations.

(b) Instead of $\Delta_t^{\epsilon} = E\{ |X_t^{\epsilon} - X_t|^{2m} \}$, we can consider the quantity

$$\widetilde{\Delta}_t^{\epsilon} = E\{ \sup_{t \in [t_0, T]} | X_t^{\epsilon} - X_t |^{2m} \}$$

as a measure of closeness between the processes X^{ϵ} and X. If we establish a sup-version of Theorem A and use some additional arguments, we can provide conditions under which $\widetilde{\Delta}_T^{\epsilon} \rightarrow 0$ as $\epsilon \rightarrow 0$ on finite fixed intervals or on intervals whose lengths tend to infinity.

(c) The results of the present paper can be used when studying stability properties of SDE's under perturbations. Another possibility is to look for the so-called expansions of the solution X^{ϵ} of the perturbed SDE (3) assuming some smoothness of the coefficients $\tilde{a}(\cdot)$ and $\tilde{b}(\cdot)$.

(d) Similar questions can be raised for more general SDE's driven by arbitrary semimartingales not just by the standard Wiener process (see Protter [6]).

Acknowledgement

One of the authors (J.S.) is grateful to Professor Tom Stroud (Queen's University, Kingston, Canada) for his attention and support. We are grateful to the anonymous referee and the associate editor for their useful comments.

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